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MEMORANDUM REPORT ARBRL-MR-03306

DRIFT VELOCITY COMPUTATIONS FOR
SHAPED-CHARGE JETS

Steven B. Segletes

September 1983

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US ARMY ARMAMENT RESEARCH AND DEVELOPMENT COMMAND
BALLISTIC RESEARCH LABORATORY
ABERDEEN PROVING GROUND, MARYLAND

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REPORT DOCUMENTATION PAGE		READ INSTRUCTIONS BEFORE COMPLETING FORM
1. REPORT NUMBER MEMORANDUM REPORT ARBRL-MR-03306	2. GOVT ACCESSION NO. AD-A133756	3. RECIPIENT'S CATALOG NUMBER
4. TITLE (and Subtitle) DRIFT VELOCITY COMPUTATIONS FOR SHAPED-CHARGE JETS		5. TYPE OF REPORT & PERIOD COVERED
7. AUTHOR(s) S. B. Segletes		6. PERFORMING ORG. REPORT NUMBER
9. PERFORMING ORGANIZATION NAME AND ADDRESS US Army Ballistic Research Laboratory ATTN: DRDAR-BLT Aberdeen Proving Ground, MD 21005		8. CONTRACT OR GRANT NUMBER(s)
11. CONTROLLING OFFICE NAME AND ADDRESS US Army Armament Research & Development Command US Army Ballistic Research Laboratory (DRDAR-BLA-S) Aberdeen Proving Ground, MD 21005		10. PROGRAM ELEMENT, PROJECT, TASK AREA & WORK UNIT NUMBERS 1L162618AH80
14. MONITORING AGENCY NAME & ADDRESS (if different from Controlling Office)		12. REPORT DATE September 1983
		13. NUMBER OF PAGES 81
		15. SECURITY CLASS. (of this report) Unclassified
		15a. DECLASSIFICATION/DOWNGRADING SCHEDULE
16. DISTRIBUTION STATEMENT (of this Report) Approved for public release; distribution unlimited.		
17. DISTRIBUTION STATEMENT (of the abstract entered in Block 20, if different from Report)		
18. SUPPLEMENTARY NOTES		
19. KEY WORDS (Continue on reverse side if necessary and identify by block number) shaped-charge jet drift distribution penetration sidewall impact drift velocity transverse velocity		
20. ABSTRACT (Continue on reverse side if necessary and identify by block number) bet/ This work provides quantitative support for the assumption that jet curvature is the major cause for nonideal jet penetration. Jet radiographs were analyzed to ascertain the component of jet velocity normal to the axis of symmetry of the charge for each jet particle. This nonideal velocity component is the drift velocity. The drift velocity distribution determined in this manner was imposed on the jet in a computer simulation of the penetration process. The computer code, referred to as PENJET, was employed to generate penetration-standoff		

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curves using the data from a specific round. This curve depicts the penetration-standoff performance for the single round as if it could be shot repeatedly at a variety of standoffs. The curve could then be compared with the actual datum for the round.

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I. INTRODUCTION

Theories to predict penetration of shaped charges have met with varying degrees of success. Except for relatively short standoffs, penetrations diverge greatly and never really approach idealized values. Inspection of jets through the use of flash radiographs reveals why idealized penetration is not a reality: because the idealized jet is not a reality. The presence of drift velocity (seen in a radiograph as jet curvature) is suggested as being the main cause for the variance in jet penetrations. Transverse (or alternately drift) velocity refers to the motion of shaped-charge jet particles in the plane perpendicular to the charge axis of the warhead (drift or transverse plane). The distance travelled by a jet particle in the transverse plane is referred to as the particle deviance.

All shaped charge jets contain drift velocity to some degree or another. The origins of drift are naturally linked to imperfections in the warhead itself. Lack of uniformity of liner thickness, nonhomogeneous explosive, as well as asymmetric liner confinements must all be considered as possible causes of drift velocity in any given round. Fissures in the explosive material and uncentered boosters can also contribute to jet performance that is less than perfect. With respect to the rounds studied in this analysis, it was also noticed that rounds with pressed explosive possessed consistently lower penetration capabilities than their cast explosive counterparts. One theory is that the high pressures (20,000 psi) required to press the explosive were capable of distorting the liner. The problems characterizing these pressed rounds were typical of the early days of pressing. Technology in the area of pressing explosives has since improved to a point where many of the problems then inherent in the rounds studied have been corrected. Some uncased pressed rounds now show penetrations of 5.5 calibers at 24 caliber standoff. Nonetheless, the pressed rounds that were studied possessed adverse qualities that manifested themselves in poor jets. Thus, with a multitude of sources for asymmetric conditions always present, it is not surprising to find the regularity with which drift velocity occurs in jets.

An algebraic analysis can be performed to evaluate numerically the drift velocities of the particles of a given jet. Preliminary and somewhat simplified computations indicate that drift velocity alone is sufficient to account for major variations in jet penetration.

II. MEASURING THE DRIFT VELOCITY OF A PARTICLE

The flash radiograph is the best tool to perform detailed analysis of the shaped charge jet in flight. It would, therefore, be appropriate to explain the nature and setup of the apparatus and how it relates to the data obtained and used in the analysis. A diagram of the apparatus is shown in Figure 1. The jet is radiographed from three different angles at three different times before it finally penetrates the target. The angular separation of each of the x-ray tubes is 45°. A sample radiograph of one of the rounds studied is shown in Figure 2. Note the fiducial markings which

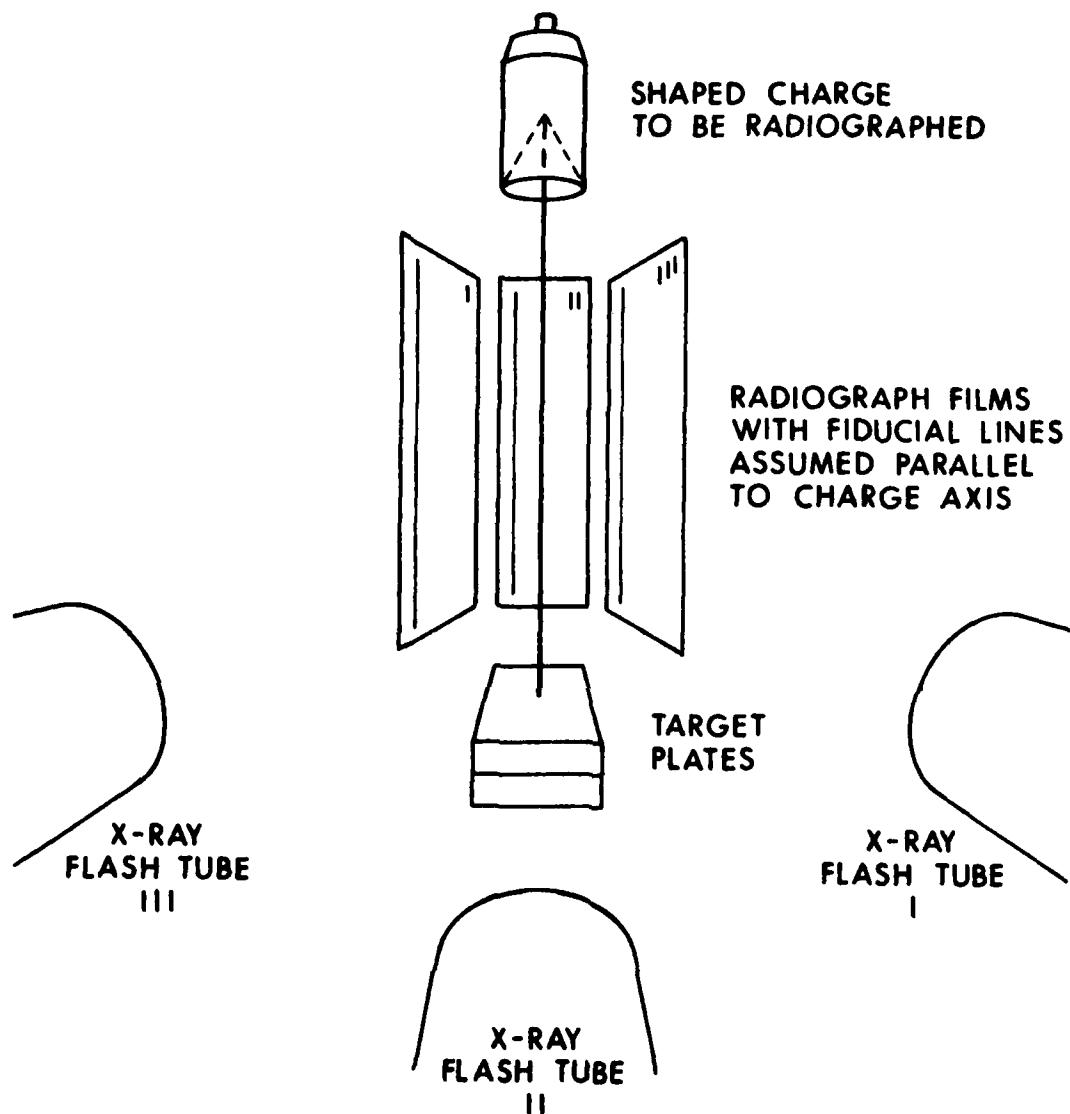


Figure 1. Radiographic Apparatus Used to Record Flight of Shaped-Charge Jet

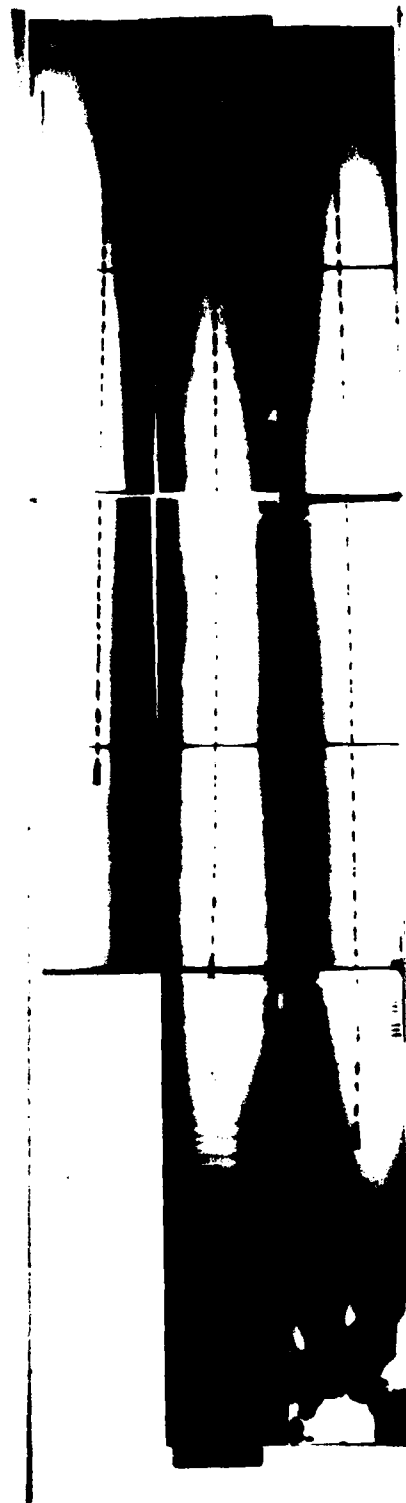


Figure 2. Sample Radiograph of a Shaped-Charge Jet in Flight

lie not on but parallel to the charge axis. Data are extracted from the radiographs by a digitizing process whereby each particle on the film is defined by six points located on the perimeter of the particle. By this means, calculations are made for each particle's length, diameter, velocity, mass, momentum, and kinetic energy using procedures described by Blische and Simmons.¹

If it is assumed that the fiducials are parallel to the charge axis of a round, the radiograph of an ideal jet would contain an image of the jet parallel to the fiducials on all exposures. This assumption is not strictly true, but shall be discussed in Section IV. With near perfect jets rarely being observed, it becomes important to measure the transverse movement of particles with respect to the charge axis. However, the location of the charge axis is, in itself, unknown. One assumes that the charge axis is parallel to the fiducials on all radiograph exposures, though it is not coincident with any one of them. It becomes necessary, therefore, to choose a workable reference. The rear of the jet has always been considered the most stable portion. For this reason, a particle near the rear of the jet that is readily definable in all three radiographs is chosen as the reference. The reference axis is an axis parallel to the warhead charge axis on which the reference particle lies. If the reference particle is not drifting in the transverse plane, the reference axis will be coincident with the charge axis.

To see how the transverse velocity of a particle might be measured, it would be convenient to imagine looking down the charge axis, as if from liner to target. Ignoring all particles of the jet except the reference particle and the particle in question, the charge axis view of the situation could be described by Figure 3. Note that the film shown in the figure is just the "edge" of the film since the view is down the charge axis. The deviance (perpendicular distance between the reference axis and the particle in question) that is seen in the radiograph is less than or at most equal to the actual deviance due to the cosine effect. Given the apparatus of Figure 1, and having exposed the three radiographs simultaneously, the reference particle and the particle in question could be precisely located with respect to each other in the transverse or drift plane. This situation is illustrated in Figure 4. The time delay that actually exists among the three exposures complicates matters little and three governing equations are readily derived:

$$V_T^* = \frac{D_1 / \cos(\theta + 45^\circ)}{t_1} \quad (1)$$

$$V_T^* = \frac{D_2 / \cos \theta}{t_2} \quad (2)$$

¹J. Blische, B. Simmons, "A Method for Reducing Data from Radiographs of Shaped-Charge Jets," BRL Report ARBRL-TR-02330, Jun 81 (AD A102770).

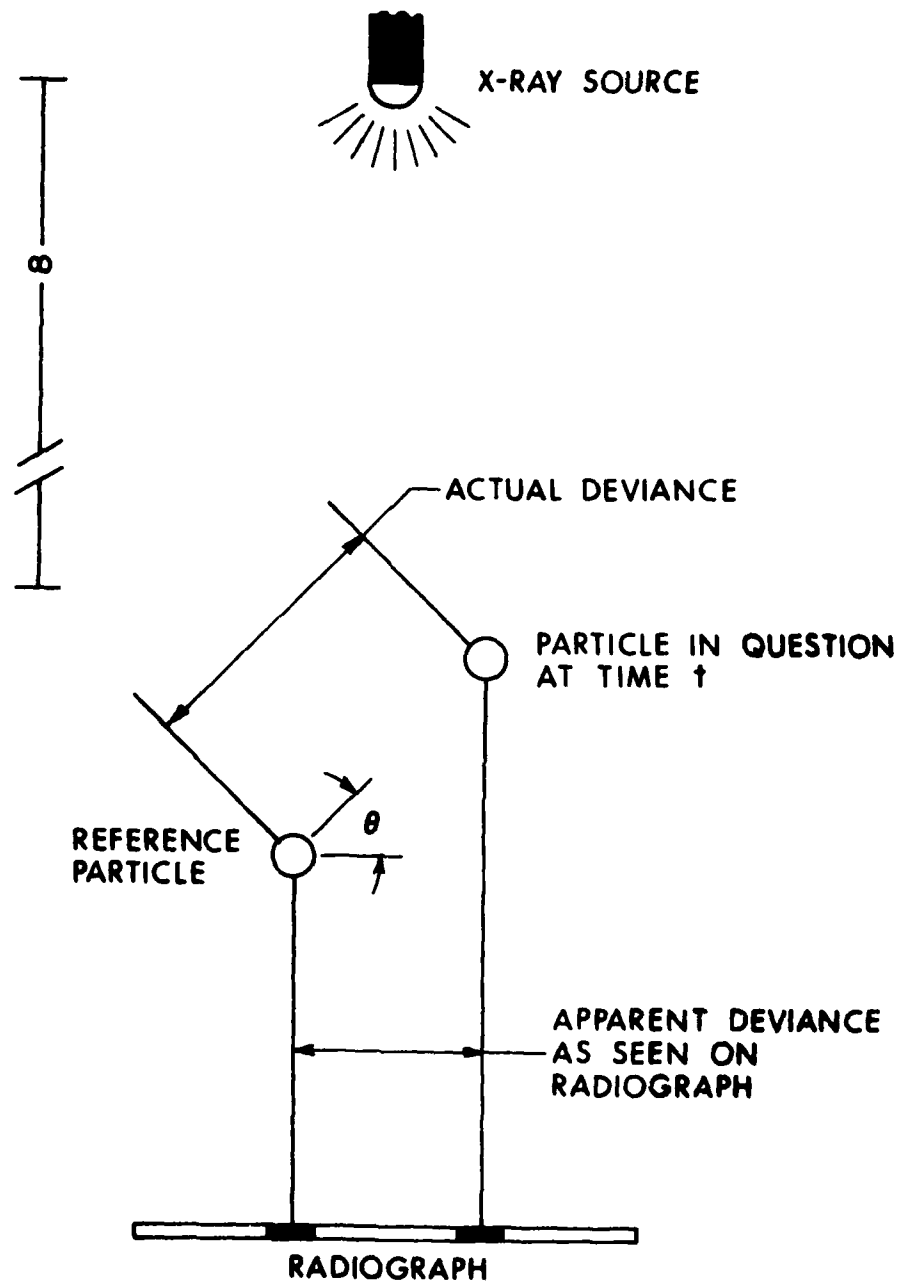


Figure 3. View of Particle Deviation as Projected onto the Transverse Plane

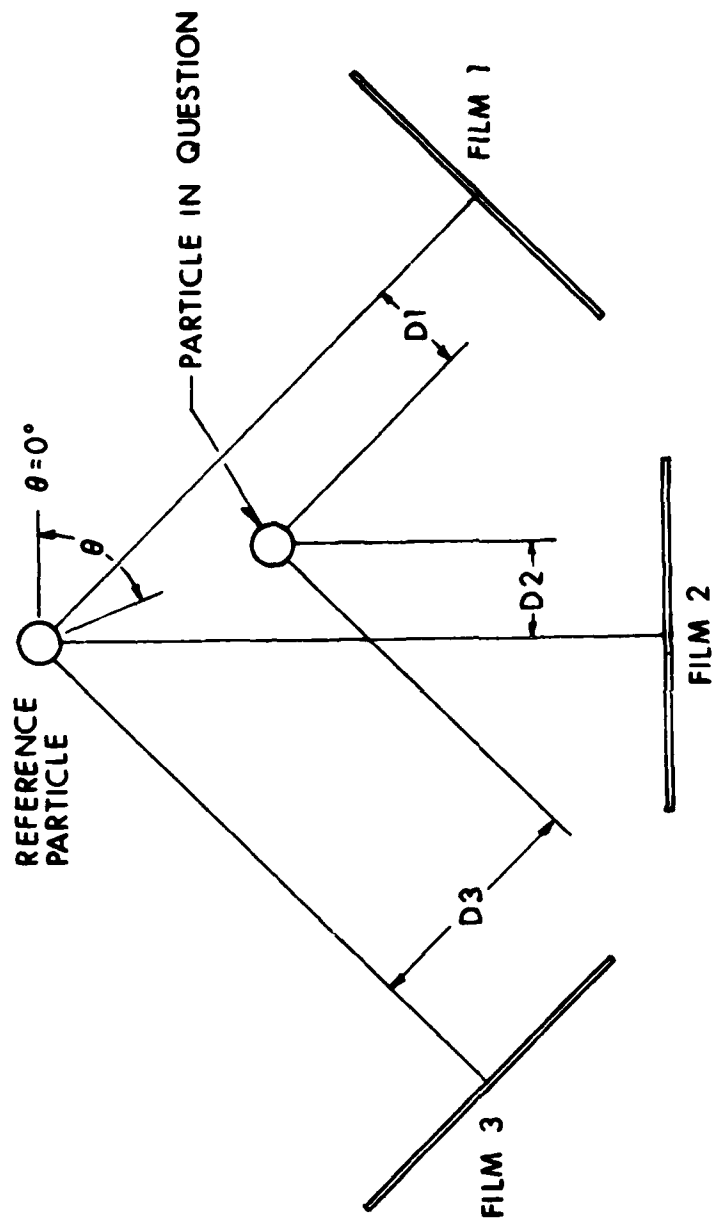


Figure 4. Locating a Particle in the Transverse Plane Through the Use of Three Radiographs

$$V_T^* = \frac{D_3 / \cos(\theta - 45^\circ)}{t_3} \quad (3)$$

where D_k - apparent deviance read off the K'th radiograph

t_k - the time between particle formation and the K'th exposure

θ - deviation angle between path of deviation and some arbitrary reference (chosen parallel to film 2)

Equations (1), (2), and (3) all contain two unknowns, namely magnitude of drift velocity (V_T^*) and angle of deviation (θ). Any two of these three equations can be solved simultaneously to yield a result for transverse velocity, both in magnitude and direction. In reality, the equations are solved in all three combinations, and a weighted average is taken depending on the particle's orientation with the radiographs. This velocity shall be called by the name of relative drift velocity since the particles are moving with respect to a moving reference (the reference particle).

III. THE DRIFT VELOCITY DISTRIBUTION FOR A GIVEN ROUND

When relative transverse velocities are evaluated for all the particles of a jet, a polar plot can be made whereby a line segment with length proportional to the magnitude of drift velocity is plotted at the appropriate deviation angle in the transverse plane. A sample velocity fan is shown in Figure 5. This relative velocity fan shows a kind of spray of the jet particles away from the reference particle in the deviation plane.

When a plot is made of the axial versus transverse particle velocity for a given round, a graph similar to Figure 6 frequently results. The data seems to be of the form:

$$V_A = KV_T^* + V_{REF} \quad (4)$$

where: V_A - axial velocity of particle

V_T^* - relative transverse velocity of particle

K - constant varying from round to round

V_{REF} - axial velocity of the reference particle

Most rounds studied were in high correlation with Equation (4).

The next quantity to determine is the drift velocity of the reference particle from which all the calculations were made. If we assume that the reference drift velocity is negligible, then the separation distance between impacts of the jet tip and the reference particle would routinely exceed

ROUND 2344

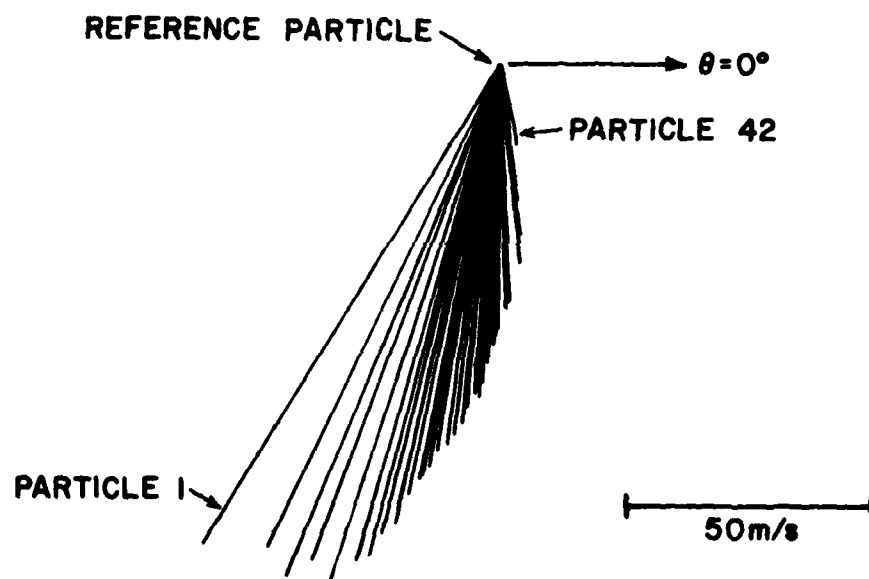


Figure 5. Relative Velocity Fan Showing the Pattern Formed by the Drifting Jet Particles

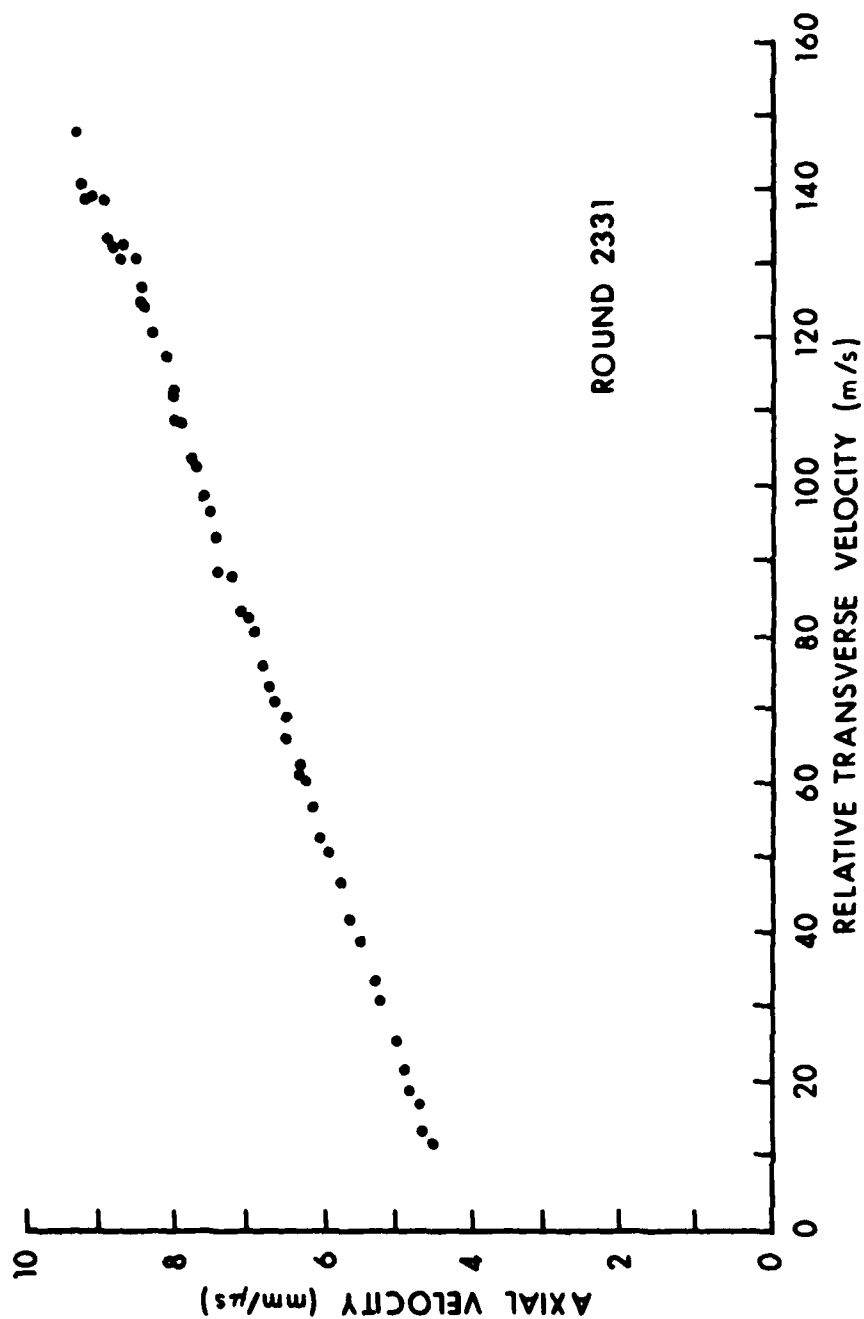


Figure 6. Typical Plot of Axial vs. Transverse Velocity for Particles of a Given Round

1/2 charge diameter (CD) at a 24 CD standoff for typical drift distributions studied. Under such conditions, the reference particle will never enter the crater created by the jet tip because of excessive drift. However, it is not uncommon for a jet to produce a single entrance hole into the target at 24 CD standoff. Thus, the reference must possess a drift velocity which will tend to bring the reference into the crater formed by the jet tip. An equation is needed which expresses that the reference particle should deviate the same distance from the charge axis as the jet tip upon reaching a target block at a given standoff. Mathematically, the assumption is:

$$D_R = D_j \quad (5)$$

where: D_R - the distance the reference has deviated upon reaching the target surface (magnitude only)
 D_j - the distance the jet tip has deviated upon reaching the target surface (magnitude only)

Equation (5) does not preclude drift among the jet particles since it pertains to magnitude only. An angular separation of the drifting particles still exists and will tend to separate the jet particles at extended standoff. For a small angular deviation among jet particles however, Equation (5) gives jets the ability to produce single crater target impacts. It is suggested solely because a compensatory reference particle drift was required and is the best which can be proposed at this time. Since drift velocity is assumed to be constant after a particle's formation, Equation (5) can be expressed as:

$$V_{T_R} t_R = V_{T_j} t_j \quad (6)$$

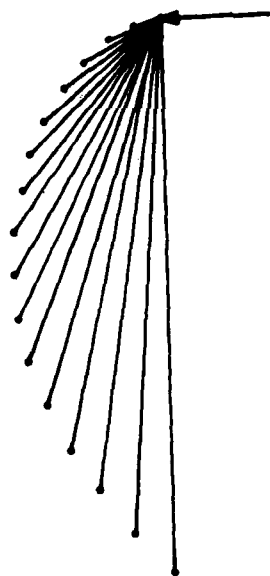
where: V_T - the absolute drift velocity of the associated particle
 t - the time required for the associated particle to reach the target surface
 R - the reference particle
 j - an arbitrary jet particle

But the time required for a particle to reach the target surface is equal to the standoff as measured from the virtual origin (Z_o) divided by the particle's axial velocity (V_A). Therefore, Equation (6) expands to:

$$V_{T_R} \frac{Z_o}{V_{A_R}} = V_{T_j} \frac{Z_o}{V_{A_j}} \quad (7)$$

Simplifying:

$$\frac{V_{A_R}}{V_{T_R}} = \frac{V_{A_j}}{V_{T_j}} \quad (8)$$



A RELATIVE VELOCITY FAN WITH SCALED REFERENCE VELOCITY ARROW INDICATING MAGNITUDE AND DIRECTION OF REFERENCE DRIFT. THE APEX OF THE FAN REPRESENTS THE REFERENCE PARTICLE FROM WHICH DRIFT MEASUREMENTS WERE ORIGINALLY TAKEN.

AN ABSOLUTE VELOCITY FAN IS CREATED BY CONNECTING THE TAIL OF THE REFERENCE VECTOR TO THE RELATIVE VELOCITY FAN. THE APEX OF THE FAN LOCATES THE CHARGE AXIS IN THE TRANSVERSE PLANE. THE REFERENCE PARTICLE NOW BECOMES JUST ANOTHER DRIFTING PARTICLE.

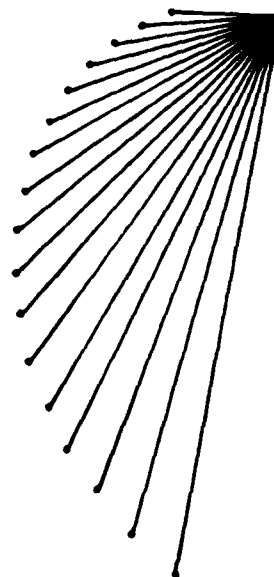


Figure 7. Transformation of Relative Fan Into Absolute Fan for a Hypothetical Drift Distribution

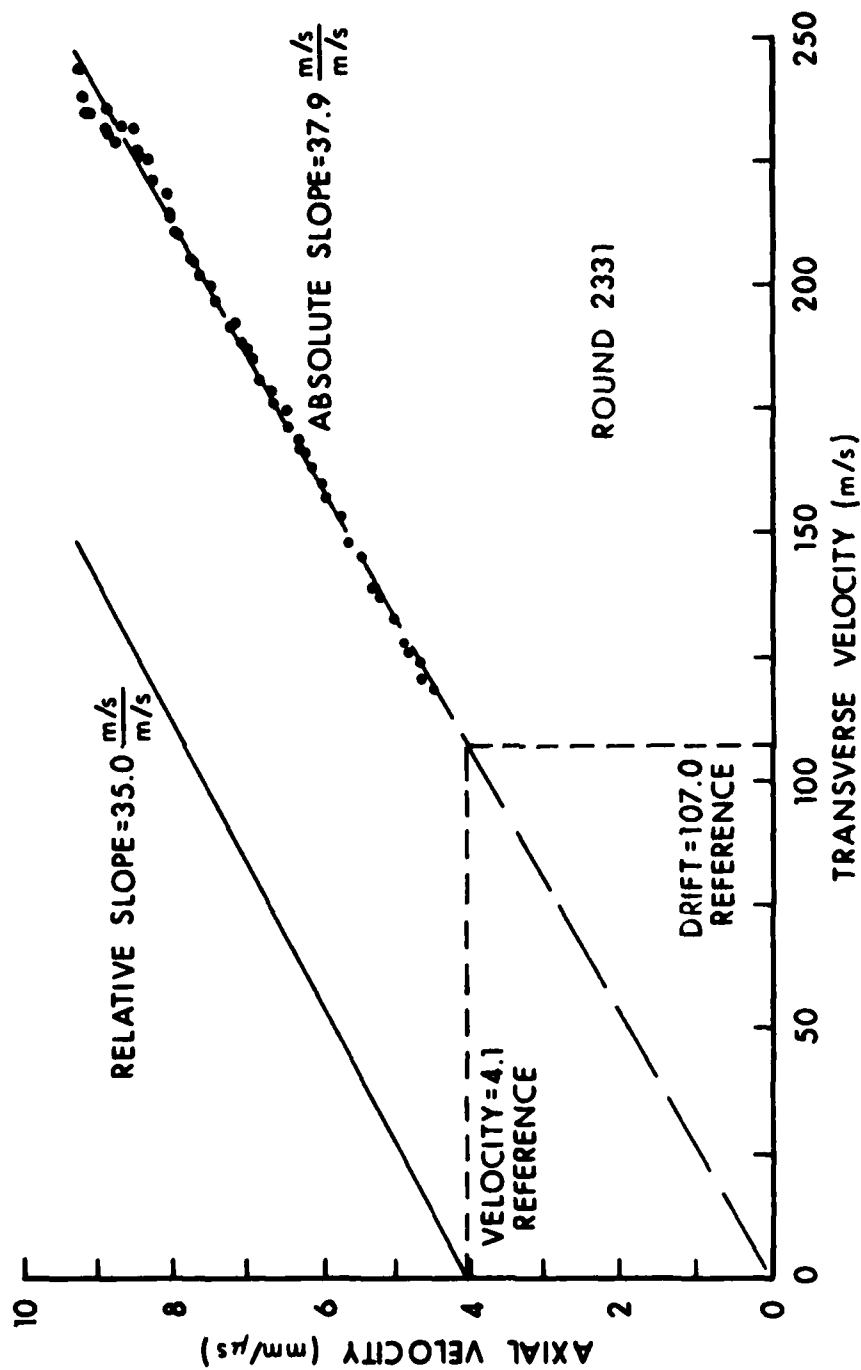


Figure 8. Effect of Transformation from Relative to Absolute Drift Velocity on the Slope of the Curve

For a crater profile in which all particles of the jet enter the target surface at the same point, Equation (5) must hold for all particles, implying:

$$\frac{V_A}{V_T} = \text{constant} \quad (9)$$

Since the reference is at the apex of the relative transverse velocity fan, the reference velocity would appear as a vector with its tip attached to the apex of the relative velocity fan. Its length corresponds to the drift velocity of the reference particle and its direction indicates the direction in which the reference is deviated. Its direction of deviation is assumed to be about the same as that of the slowest digitized particle. The absolute velocities of the balance of the particles are calculated by vectorially adding the reference velocity to the relative velocities of the respective particles. Thus, an absolute drift velocity fan is created by drawing the segments from the tail of the reference vector to the tips of the relative drift velocity fan. (See Figure 7.) Graphically, Equation 9 implies a linear V_A vs. V_T plot which intercepts the origin. With the proper vectorial addition of the reference drift as just described, relative drift distributions like Figure 6 readily map into absolute drift distributions which satisfy Equation 9. (See Figure 8.)

There is no way to know whether this transformation into absolute velocities is the correct one. However, it is believed that the directional assumption of the reference is a safe one, and that the Equation (5) assumption is as good as any considering the need for the drift correction of particles. In addition, the results obtained from the transformation are supported by the experimental data. Note that Equation (9) does not guarantee that all particles will only form a one-crater hole profile even though all the particles should deviate the same distance from the charge axis. A large angular spread in a jet (characterized by a wide absolute velocity fan) may produce enough relative deviation among the particles so as to cause multiple impacts at long standoff. Thus, the deviating factors in jets are both the angular spread and the magnitudes of the drift velocities.

IV. THE QUESTION OF TILT

A possible and potentially fatal flaw in the previous discussion concerns the presence of tilt. The word tilt is taken to mean a charge axis that is not parallel to the radiograph fiducials. An original assumption was that there was no tilt, or that charge axis and fiducials were aligned. It is obvious that some sort of tilt must exist, however small. The answer as to when tilt becomes critical is a bit disconcerting. Depending on the jet, a one-degree tilt of the charge axis off the fiducials could induce an apparent drift velocity of 160 m/sec in the jet tip as measured with respect to the fiducials. It is likely that uncertainties of this magnitude exist regularly.

The consequence of tilt on a jet is an apparent drift velocity distribution, which is also governed by Equation (9), which is the assumed relation for actual drift velocity. One may rightfully wonder whether all of the drift velocities read off the radiographs are due completely to tilt, and whether graphs like that of Figure 6 are a consequence of tilt alone. In fact, what is read off the radiographs is the superposition of both tilt and real drift velocity distributions. However, it should be noted that an important aspect of tilt is that it cannot produce a velocity fan. All of the particles of a jet under the influence of tilt exclusively would appear to drift in a uniform direction. Yet, Equation (9) has been seen to hold for jets with velocity fans spanning in excess of 90° .

The concern here is whether a nonlinear drift distribution could be "tilted" into appearing linear on the radiographs, and thus explain the data without having the actual drift obey Equation (9). Such could only be the case if the magnitude of actual drift were small in relation to tilted drift. If this were the case, however, the angular spans of the velocity distributions would also be quite small, which is usually not the case. Thus, it appears that the actual drift velocities are at least of the same order of magnitude as the apparent drift velocities due to tilt. Thus, it seems that the actual drift distributions, though probably different from the measured drift distributions, can still be adequately described by Equation (9). This belief is based only on the rounds studied. No explanation is offered as to why there might be a relation between the axial and transverse velocities of jet particles. Also, it is uncertain as to whether this may be a general phenomena or characteristic only of the rounds studied.

Support for Equation (9) was found in work done by Aerojet General Corporation.² Aerojet was interested in the crater formation of individually impacting particles. A shaped charge supplied the particles for their experiment. The method employed for dispersion of the particles was asymmetric initiation. Considering that it was desired to induce much drift velocity in the particles, the effects of tilt could be neglected in this experiment. The target plate used in their experiment is shown in Figure 9. The particle impacts generally lie at a constant distance from the slug impact. Assuming the slug flight path to be coincident with the charge axis, Equation 9 is the governing equation.

A penetration code was created to predict penetrations for nonideal shaped-charge jets. Unlike the DSM model,³ which can express penetration in a single equation, the model developed analyzes the jet piecemeal. It is described in detail in the following section.

² K. N. Kreyenhagen, J. E. Ferguson, R. R. Randall, J. P. Joyce, "Special Explosive Projectors," *Proc. 6th Symposium Hypervelocity Impact*, Vol. I (Held Cleveland, Ohio, April 30-May 2, 1963).

³ R. DiPersio, J. Simon, A. Merendino, "Penetration of Shaped-Charge Jets Into Metallic Targets," *Ballistic Research Laboratory Report No. 1296*, September, 1965 (AD 476717).

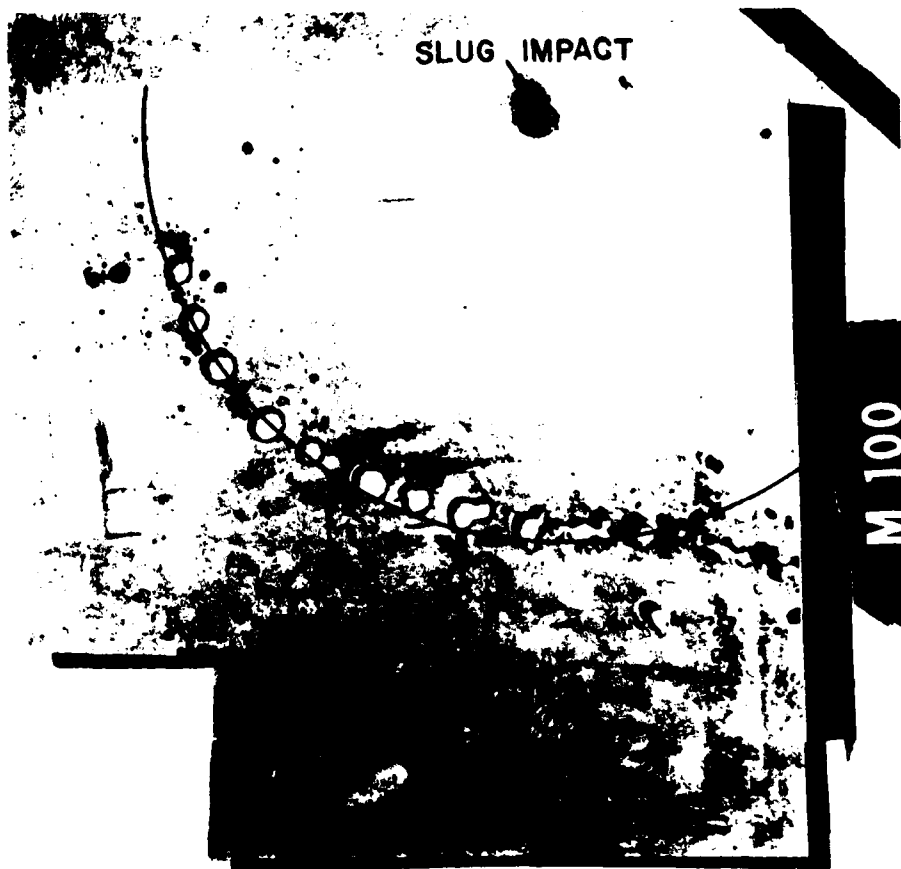


Figure 9. Target Plate for Asymmetrically Initiated Shaped Charge (Reference 2)

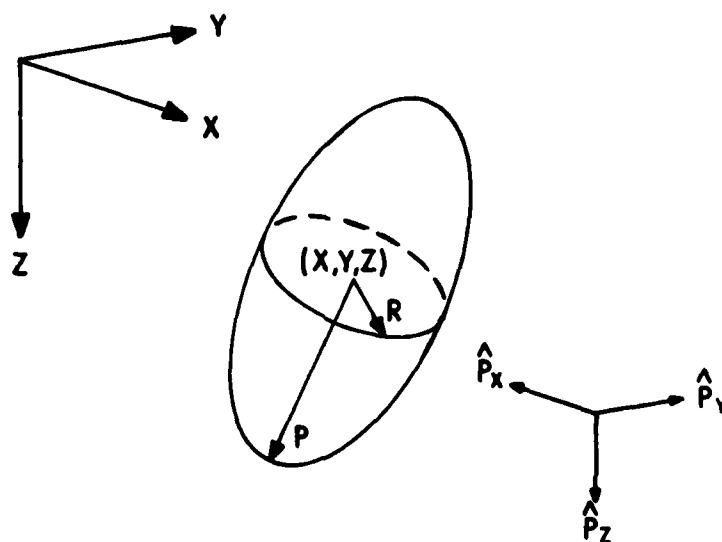
V. THE CREATION OF A PENETRATION CODE TO ANALYZE DRIFT

The purpose of analyzing drift velocity distributions of shaped-charge jets is to see how these distributions affect penetration. The dynamics of penetration is not, however, a thoroughly understood subject. Many theories exist in attempts to describe the mechanical properties of matter. Classical, empirical and even wave mechanical concepts are used in pursuit of the subject. Consequently, rather than adding to the spectrum of theories that already exist, the author chose to draw upon information already available on the subject whenever possible and to combine views in a seemingly compatible manner.

One shaped-charge penetration model that has received much attention in recent years was formulated by DiPersio, Simon, and Merendino (the DSM model).³ This model does a good job at predicting penetration of shaped charges even at extended standoffs. The model accomplished this however, by implying that there exists a cutoff penetration velocity, a minimum penetration velocity below which a jet cannot penetrate. Unfortunately, the DSM cutoff penetration velocity implies that jet material traveling below certain hypervelocity speeds is incapable of penetrating target material. Substantial quantities of jet material are thus eliminated from the penetration process. The DSM model is empirically based and really attempts only to describe the data at hand and does not enter an intense theoretical discussion to describe the origins of jet cutoff other than citing that it exists experimentally. The work presented in this report suggests that the drift of particles off axis is largely responsible for reduced jet penetration, whereby drifting particles strike the target not at the bottom of the hole but somewhere up the hole profile. This commencement of sidewall impacts frequently seems to constitute a cutoff in jet penetration, since the overall depth of penetration is not altered by a sidewall impact.

A method, therefore, had to be devised which would allow penetration to be a function of drift velocity. Three-dimensional finite difference codes were not considered due to a coupling of availability and cost. Yet, a three-dimensional description of penetration was required since the source of altered penetration is the spread of the particles off of the one-dimensional axis. Consequently, all aspects of particle penetration would have to be considered: size, shape, and orientation of crater formation in a three-dimensional target block.

The code developed in the present work (referred to as PENJET) includes many simplifications and assumptions. The first assumption concerns the shape of crater formed by an impinging jet particle. The chosen shape was that of an ellipsoid, with the ellipsoid being characterized by a position in three-dimensional space, dimensions along the axes, as well as an orientation in space (as shown in Figure 10). The position of particle impact is determined by using the velocity fan of a round to trace the particles' flight through three-dimensional space and time in order to locate the particle position with respect to the instantaneous hole profile (formed by earlier particles of the jet). If a particle is found to strike any point within the target, PENJET then proceeds through a series of calculations to determine the remainder of crater parameters. Once the crater formed by a particle is effectively evaluated, PENJET proceeds through subsequent particles of the jet until impacts for all digitized particles have been evaluated. Thus,



(X, Y, Z) LOCATION OF IMPACT IN 3-D SPACE
 P MAGNITUDE OF PARTICLE PENETRATION
 R MAGNITUDE OF CRATER RADIUS
 $\hat{P}_x, \hat{P}_y, \hat{P}_z$ PENETRATION AXIS
 (A UNIT VECTOR IN 3-D SPACE)

Figure 10. Parameters Generated by PENJET to Describe Crater Formed by Impinging Jet Particle

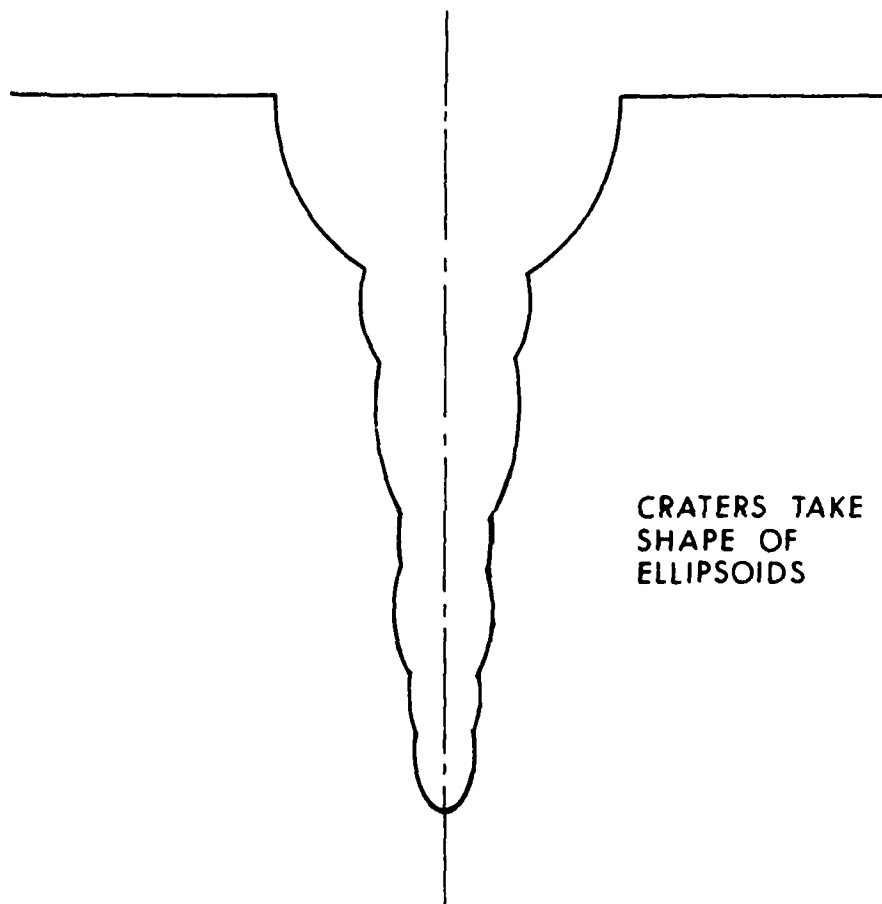


Figure 11. PENJET Hole Profile Results from Ellipsoid Superposition

a hole profile for a complete jet is just the superposition of ellipsoids from each of the particle impacts. A two-dimensional representation of a three-dimensional hole profile is shown in Figure 11 for a fictitious nondrifting jet. When evaluating the crater formed by an impinging jet particle, crater volume is a constraining factor. That is to say, crater volume is calculated independently of other crater parameters and is assumed to be a function of particle kinetic energy alone. Subsequent calculations for actual crater dimensions are, therefore, constrained by the allowable volumetric target displacement a given particle can produce. The function governing the volume energy relationship is a linear one, similar to that adopted by DiPersio, Simon, and Merendino.³ The distinction arises in the fact that the DSM model expresses its proportionality constant as a relationship between total jet kinetic energy and total hole volume. PENJET, on the other hand, bases its constant on the energy of the individual particle. It is not surprising to find, therefore, that the DSM constant differs from its PENJET counterpart. Consider also, that the energy constant of PENJET does not strictly define an exact amount of displaced target material. As mentioned previously, the hole profile of PENJET is the superposition of ellipsoid craters corresponding to each of the jet particles. The orientation of these ellipsoids will greatly affect the spacial overlap of the ellipsoids, and thus the overall volumetric target displacement of a given jet. Because of this spacial overlap of particle craters, the overall crater volume as calculated by PENJET is not the summation of particle energies divided by the energy constant. Estimation of the constant was gotten empirically through examination of target blocks as well as hole profile data.

The crater shape of an individual impact is prescribed (as an ellipsoid) and the particle's crater volume is proportional to its own kinetic energy. Therefore, knowledge of a particle's penetration capability will allow the particle's hole radius to be computed and vice versa. In PENJET, depth of penetration is computed, with crater radius being subsequently determined. The model of particle penetration used in the code is based on the work of Fitzgerald,⁴ who draws on the work of Kineke⁵ to imply that crater formation occurs in a direction perpendicular to the stricken surface. Fitzgerald calls this phenomenon the wave refraction theory and shows that it applied to a single particle striking a flat plate at some obliquity (Figure 12). PENJET uses the concept in a more general sense by applying the wave refraction theory to local obliquities as well. Local obliquity is defined as the angle between the particle flight axis and the normal to the impact plane. Unless the impinging particle strikes exactly at the bottom of the hole profile, the refracted penetration wave will not be normal to the original target surface (Figure 13). In the strict sense, the refracted wave theory should only be applicable in the regime of hydrodynamic penetration. A large majority of particle impacts studied did not fit the qualification for hydrodynamic penetration. For these cases, the direction of penetration would be somewhere between the particle flight axis and the refracted wave axis. It was later noted that the penetration of a particle typically varied less than 5% because of the refracted wave assumption.

⁴E. Fitzgerald, *Particle Waves and Deformation in Crystalline Solids*, Interscience Publishers, New York, 1966.

⁵J. Kineke, Jr. "An Experimental Study of Crater Formation in Metallic Targets," *Proc. 4th Symposium Hypervelocity Impact*, Vol. I (Held Eglin AFB, Florida, April 26-28, 1960).

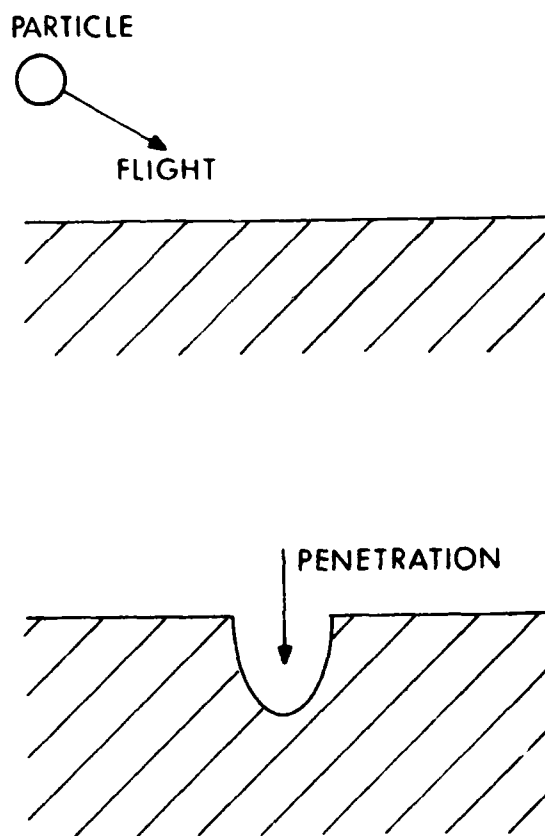


Figure 12. Crater Formation According to Refracted Particle Wave Theory

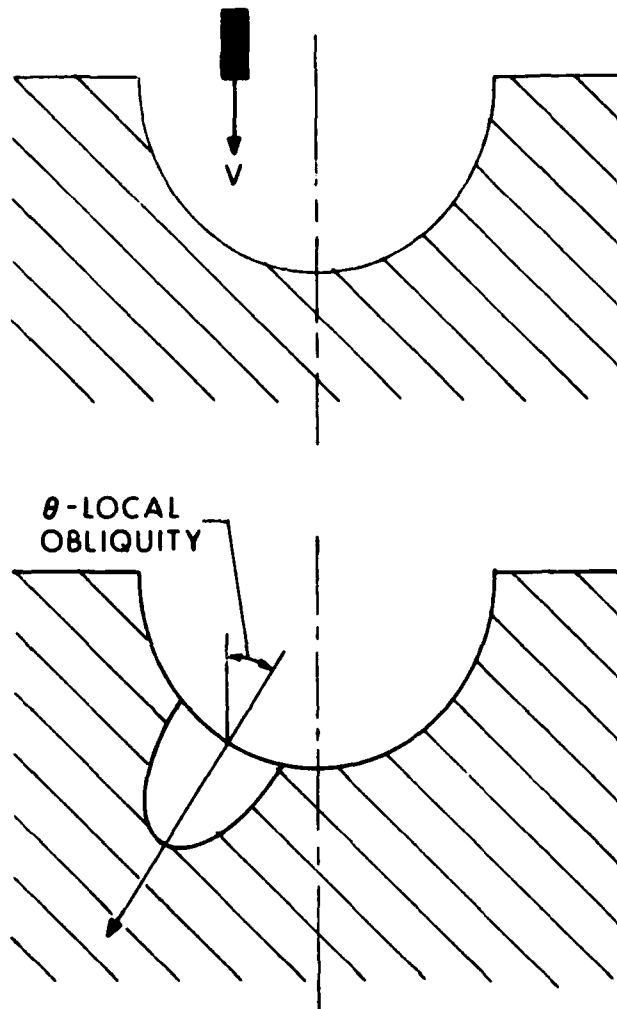


Figure 13. Refracted Particle Wave Theory Applied to Local Obliquity

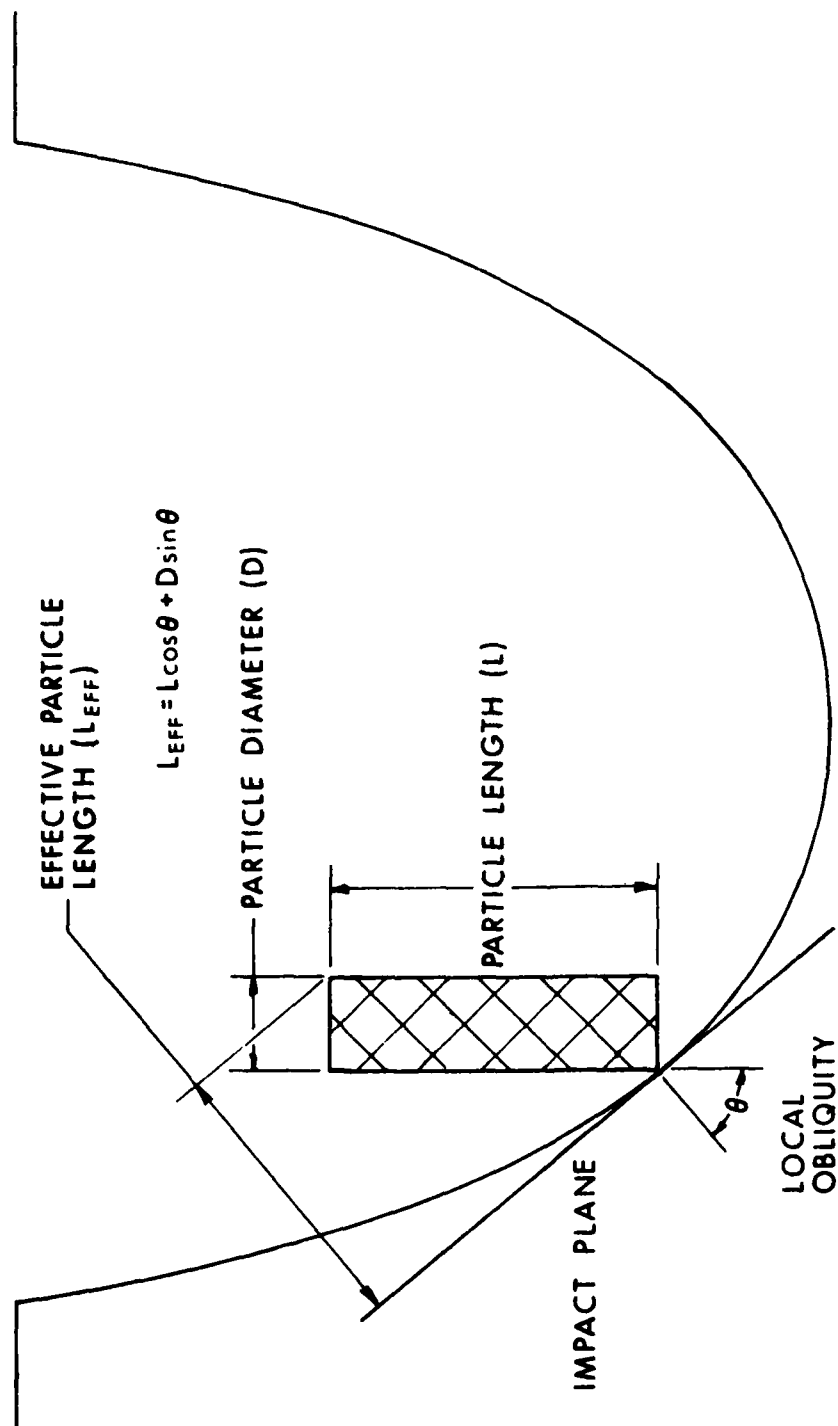


Figure 14. Concept of Effective Length as Seen by Impact Plane

Jet penetration in the continuous mode is quite a different matter from in the particulate mode. Modifications, therefore, had to be included to describe continuous jet penetration. The criterion on which to base mode of penetration is that of breakup time. PENJET uses a linear breakup equation whereby the jet tip particulates at some t_{\min} and subsequent particle breakup occurs at regular Δt intervals. Estimates of jet particulation are obtained through examination of the radiographs for the three different exposures. These estimates are then used in a least squares fit to evaluate t_{\min} and Δt . Based on the estimated breakup time, it can be established whether a given particle will be penetrating continuously or particulate at a given standoff. If continuous, the length of the particle is re-evaluated such that:

$$l_{\text{CONTINUOUS}} = \frac{t_{\text{IMPACT}}}{t_{\text{BREAK}}} l_{\text{PARTICULATED}} \quad (10)$$

Also, it is assumed that jet penetrating in the continuous mode does not experience the refracted particle wave of Figure 12. This must, of course, be the case since the continuous particle is essentially beginning its penetration exactly when and where the previous particle leaves off.

PENJET also considers particle orientation with respect to the stricken crater surface. The assumed shape of a shaped-charge particle is cylindrical, being characterized by a length and diameter. As local obliquity is increased, the length of the particle as seen by the impact point is effectively changed. (See Figure 14.) The extreme case is when local obliquity approaches 90° . In this case, the effective length of the particle approaches the diameter of the particle. Hydrodynamically, the penetration of a particle is proportional to the particle length. Thus, the effective particle length replaces the particle length in penetration calculations.

The presence of obliquity raises more questions. The larger the obliquity, the smaller the component of particle velocity perpendicular to the stricken surface. Once in the hypervelocity regime, large changes in particle velocity produce relatively small changes in penetration. If particle velocity is too small, however, the strength of the target and jet play a major role in reducing the penetration of that particle. Besides the effective particle length that the target "sees", it is believed that an effective velocity also plays a role in determining depth of penetration. Whereas penetration is assumed to be directly proportional to effective length, the role that effective velocity plays is not nearly so profound. As long as the effective velocity is in the hypervelocity regime (arbitrarily chosen as any velocity above $2\text{mm}/\mu\text{sec}$), penetration is assumed greater than the effective particle length for copper jet into an RHA target. As effective velocity increases, penetration is assumed to asymptotically approach the hydrodynamically idealized value. That is:

$$\Delta P \rightarrow L_{\text{eff}}/\gamma \text{ where } \gamma = \sqrt{\rho_T/\rho_j}.$$

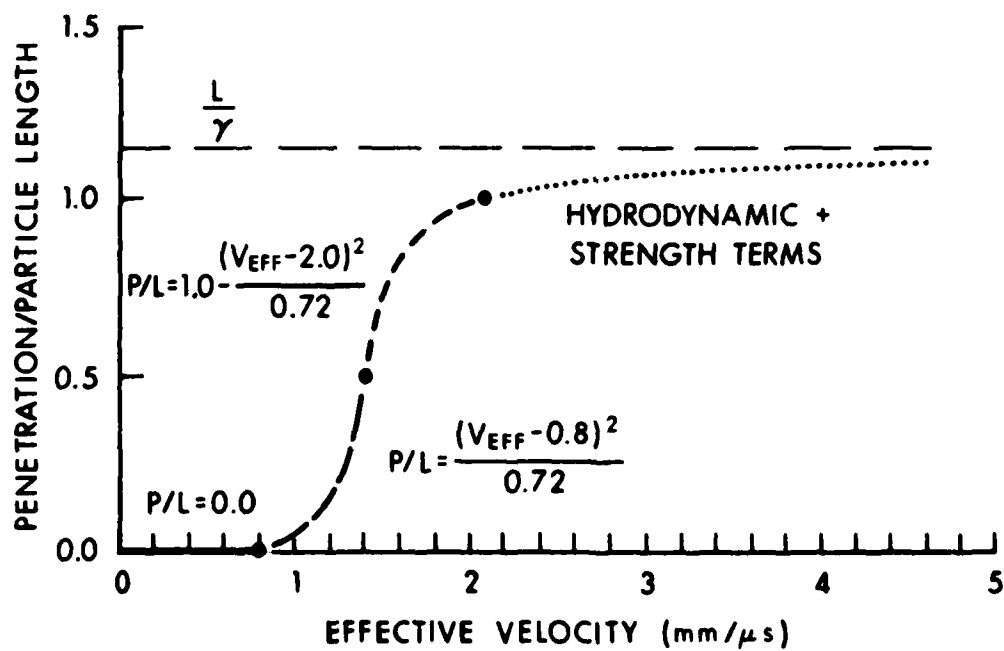
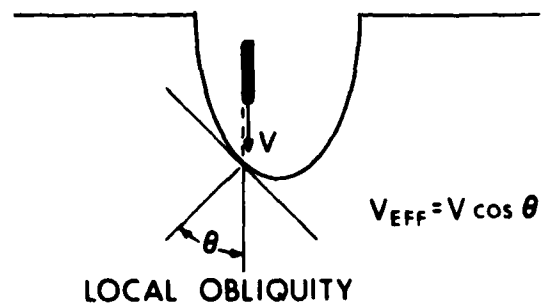


Figure 15. The Role of Effective Velocity in Penetration Calculations Involving Copper Jets Against Steel Targets

The asymptotic approach is generated using the hydrodynamic equations, including strength terms. For effective velocities below $2\text{mm}/\mu\text{sec}$, penetration rapidly decreases in a manner expressed in Figure 15 so that no penetration is induced when the effective particle velocity is lowered beyond $.8\text{mm}/\mu\text{sec}$. Data for the effective velocity curve in the low velocity regime was gathered from work by Weirauch.⁶ Impacts of this nature may be the source of particle residue shown by Simon and DiPersio⁷ to accumulate at the bottom of the hole profile. PENJET, however, does not consider this type of interference.

VI. APPLYING THE CODE TO DATA

Given a set of digitized radiographic data, PENJET will go through all of the mentioned calculations for each digitized particle of the jet to eventually create a three-dimensional hole profile for a given round at a given standoff. If the code is told to use the data of a given round to perform penetration calculations at a variety of standoffs, a theoretical penetration standoff curve can be generated for each round. This curve is, of course, only valid for a particular round. Thus, each generated curve is as individual as the data from which it was generated, and does not, in general, pertain to all rounds of the same type.

The code has both its good and bad aspects. A major consequence is that tilt has no direct effect on the code's ability to calculate penetration. This happens because tilt does not misalign the particles. Rather, it only redefines the charge axis. The particles still drift with respect to each other in the same fashion despite measurements including tilt. Much emphasis was previously placed on tilt for the reason that it would be advantageous to obtain drift distributions unbiased by tilt. In terms of the penetration calculations, though, tilt of and by itself has no effect.

Tilt does, however, indirectly affect the situation in terms of reference drift calculations. As relative drift velocity (as calculated from the radiographs) is the superposition of tilted and actual drift, abnormal relative drift distributions can occasionally result depending on how the two components of measured drift interact with each other. Correspondingly, abnormal relative drift distributions lead to abnormal extrapolations of reference drift, thus making the calculated reference drift suspect. After all, the calculation of reference drift is only an educated approximation: Direction of the drift is based on jet particle orientation, while magnitude of drift velocity is based on Equation (9). Therefore, unlike relative drift computations, which are as good as accuracy of measurement, reference drift is only as good as the governing assumptions. Thus, the poorer the approximation at reference drift, the poorer the quality of penetration calculations.

⁶ G. Weirauch, "The Behavior of Copper Pins Upon Impacting Various Materials With Velocities Between 50 m/s and 1650 m/s," (Doctoral Thesis) University of Karlsruhe, 1971.

⁷ J. Simon, R. DiPersio, "Experimental Verification of Standoff Effects on Shaped-Charge Jet Cutoff in Solid Targets," Ballistic Research Laboratory Memorandum No. 1976, May 1969 (AD 854396).

Another drawback to the code results from the radiographs themselves. The code can only analyze particles that have been digitized. In most radiographs some portion of the jet that is visible has not yet particulated. Attempting to separate the continuous jet into digitizable portions has met with some success, though likely at the cost of accuracy of particle length, since the continuous particles are still stretching when digitized. Also, most radiographs don't even contain the complete jet, for there is still jet matter that hasn't reached the plane of the film by the time it is exposed. Thus, there is no way to digitize these particles. For the rounds studied, it is estimated that over 40% of the jet was routinely nondigitizable.

To what extent the ability to digitize particles affects the penetration code predictions depends on the quality of the jet. In fact, the poorer the jet is, the less predictions are affected. If the undigitized particles drift far enough off course from the front portion of the jet, they will never make it to the bottom of the hole profile to add to penetration. In this situation, calculations based only on the particles digitized should produce results compatible with calculations involving all particles. Since distance of deviation is directly proportional to time in flight which is directly proportional to standoff, longer standoffs will also decrease the penetration effectiveness of the rear of the jet. Thus, penetration predictions at short standoff are subject to unavoidable error since the undigitized jet portion is still capable of increasing jet penetration. The fact that the penetration effectiveness of the jet rear is predicted to decrease with increasing standoff correlates with the DSM concept of minimum jet velocity increasing with standoff.³

Having kept these reservations in mind, the results of the code based on data from four rounds are illustrated in Figures 16-19. The only factor governing the choice of these particular four rounds was penetration; that is, rounds were chosen with varying penetrations ranging from poor to good. The code was run at standoff increments out to 38 CD. The actual rounds were fired at 24 CD standoff to allow radiography, and the actual penetrations achieved are shown on each of the graphs as a single datum. Steplike variations in the code's predictions result because of the discreteness of the impacting particles: Either a particle made it to the bottom of a hole profile or it did not. Major fluctuations in the code predictions can be taken to imply uncertainty in the jet performance. Uncertainty as to whether two or three given particles make it to the bottom of the hole profile can create penetration variations in excess of 1/2 CD. A curve was drawn to fit the code's predictions on each of the rounds, though a band width of uncertainty should be realized as well. This curve is PENJET's prediction of penetration, assuming that one could fire the exact same round over and over again at various standoffs. Since we know that the predictions at short standoff are inaccurate due to nondigitizable jet, the code's predictions on Figures 16-19 are generally unreliable below 10 CD. The curve, however, should be valid at the standoff of 24 CD at which the rounds were actually fired.

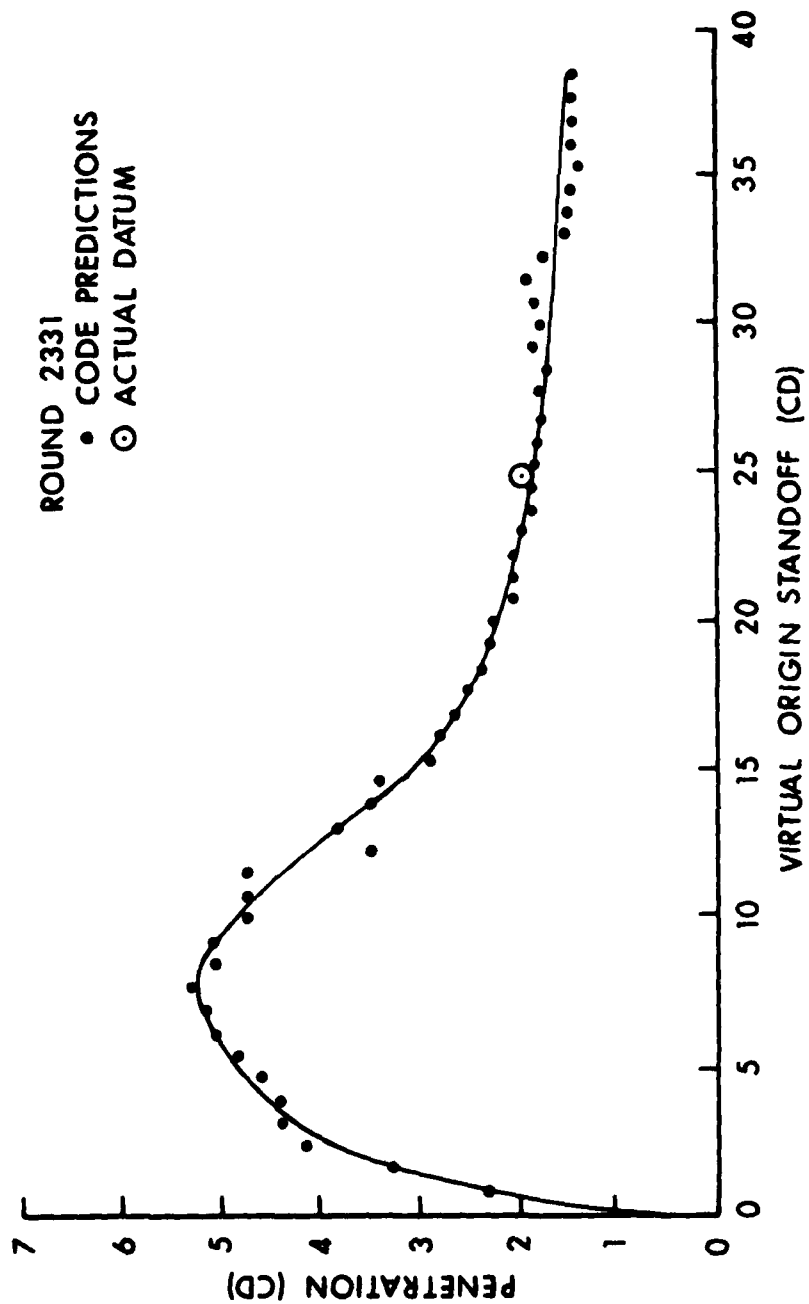


Figure 16. Analysis of Round 2331

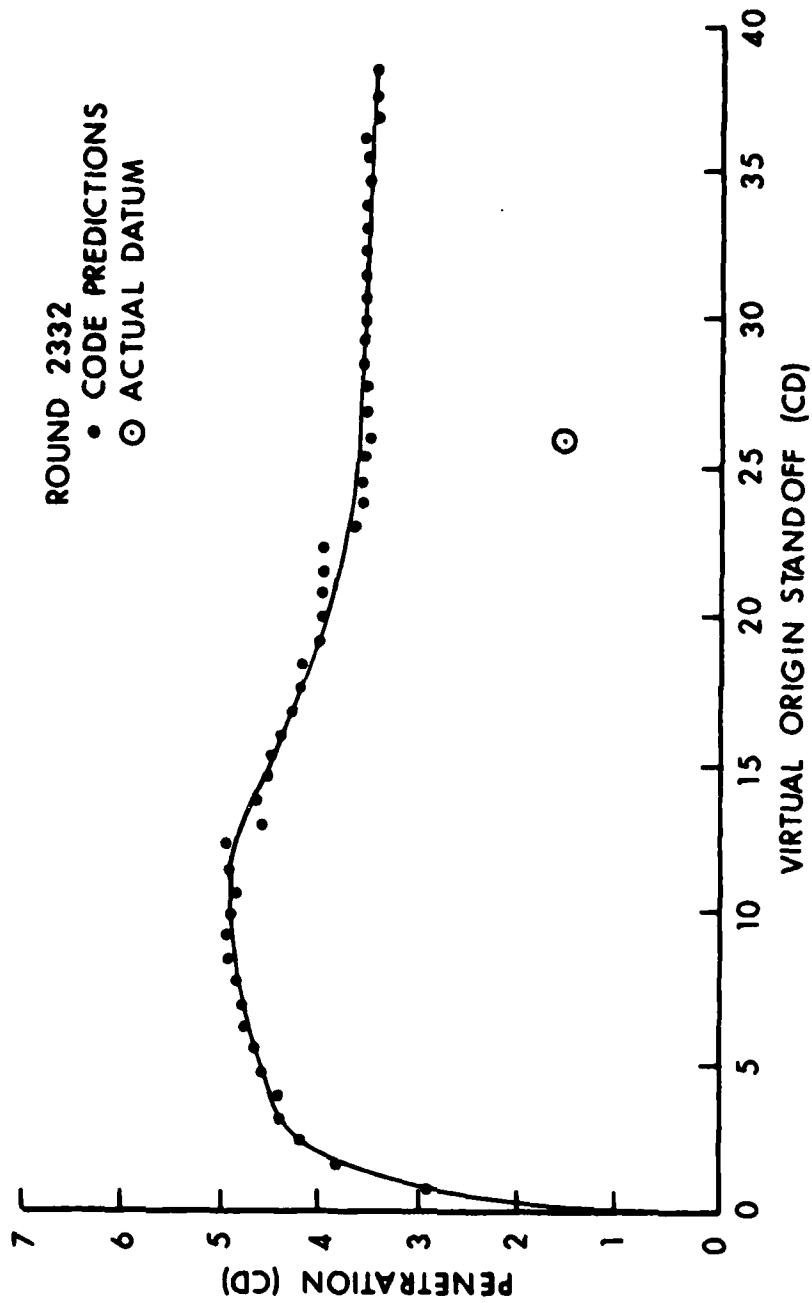


Figure 17. Analysis of Round 2332

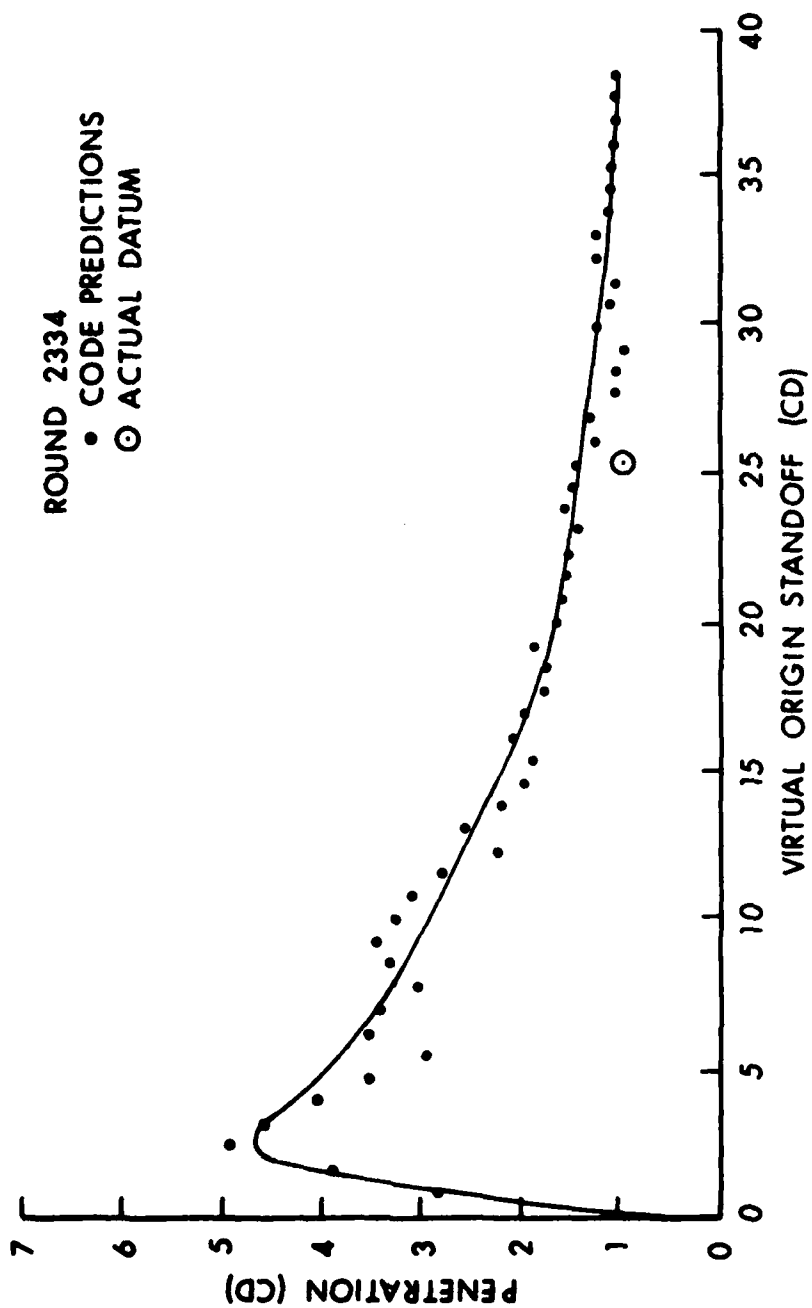


Figure 18. Analysis of Round 2334

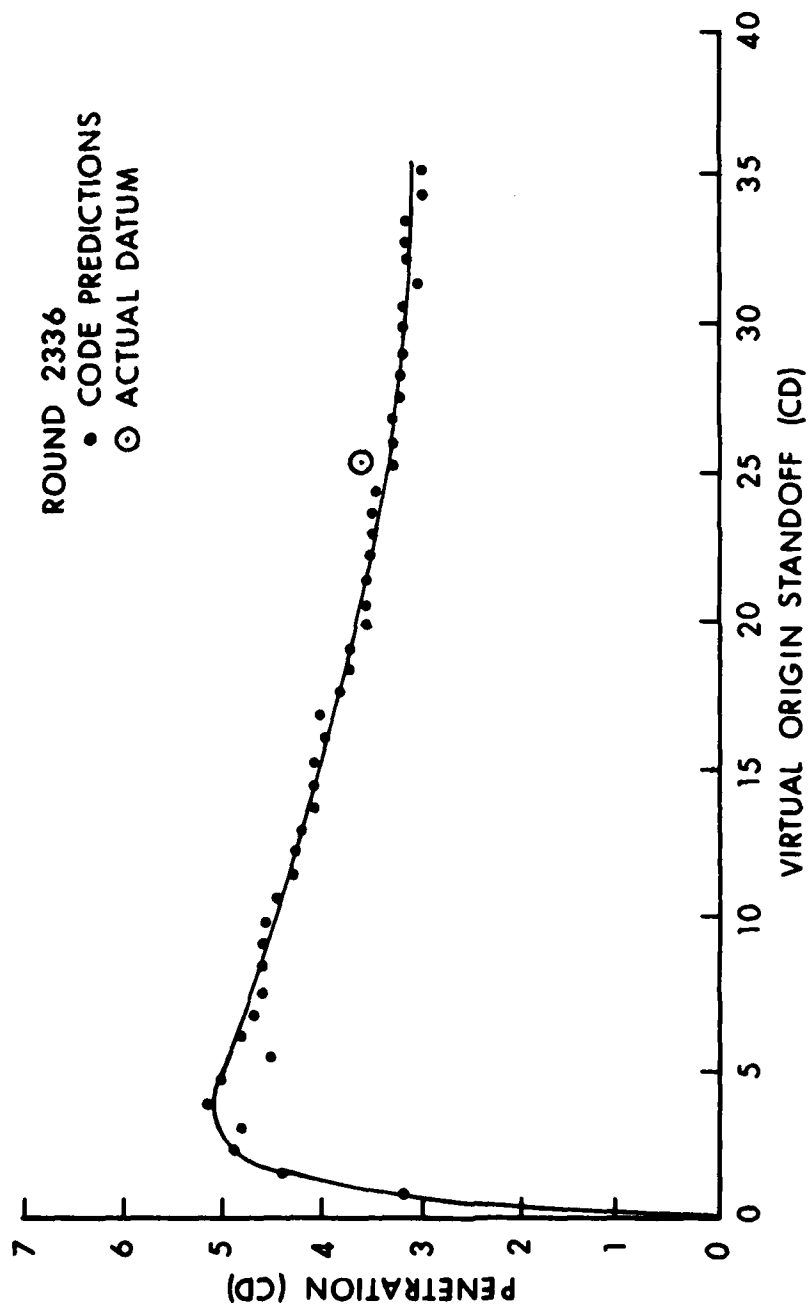


Figure 19. Analysis of Round 2336

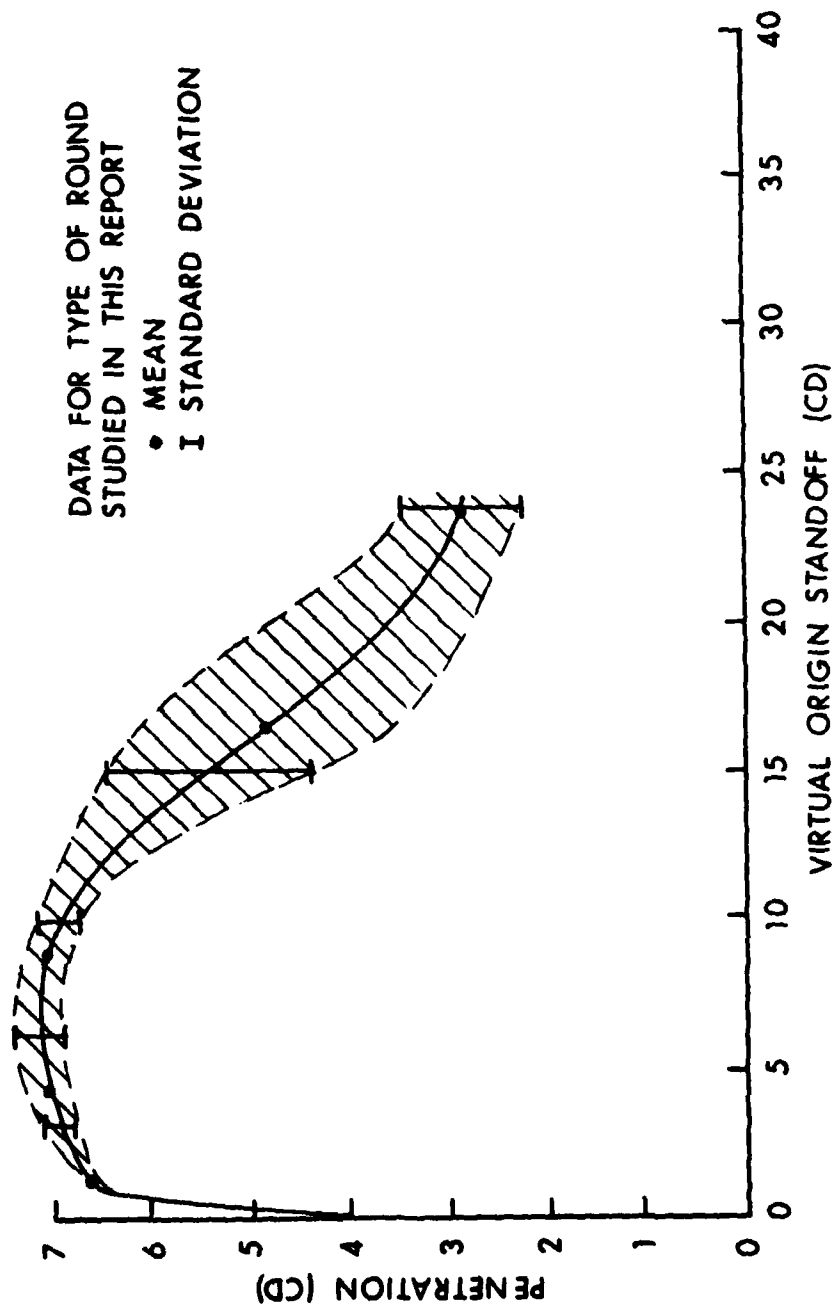


Figure 20. Routine Performance for Type of Round Analyzed

All of the codes's predictions were in reasonable agreement with experiment except for Round 2332. Upon studying this round for possible causes of discrepancy, it was noted that the round uniquely had what shall be called a switchback angular spiral. This falls into the category of abnormal relative drift distributions. Most jets contain drift distributions such that the jet particles spiral in only one direction (either clockwise or counterclockwise) throughout the jet. Figure 5, for example, shows a jet spiralling counterclockwise from jet tip to tail. A switchback angular spiral is one in which one part of the jet spirals in one direction, while the subsequent portion of the jet spirals in the opposite direction. Switchback is thought to be a product of abnormal tilt/drift combinations such that subsequent computation of reference drift might be quite inaccurate. As set forth before, an inaccurate reference extrapolation leads to inaccurate penetration calculations.

Figure 20 shows how the actual weapon routinely performs. In order to test whether PENJET could predict the average penetration at standoffs near the peak performance of the round, a radiograph containing a complete and fully articulated jet would be required for digitization. As mentioned previously, no available radiographs fulfilled these requirements. A radiograph was obtained, however, of a complete jet which was not fully articulated. Though not completely digitizable, the presence of a complete jet greatly facilitated estimation of many jet parameters (e.g. velocity of jet rear, drift with respect to jet rear, and total jet length). By this means, the rear of the jet which was not digitizable could be simulated into particles based on analysis of the radiograph containing the continuous jet rear.

Figure 21 shows the PENJET prediction of the round that includes the jet rear simulation. The actual datum as well as optimum performance for this type of round are included in order to put PENJET calculations in perspective.

The importance of a complete jet at short standoff is shown in Figure 22. The penetration standoff curve for Round 2331 is shown here. Superimposed upon it are curves showing the amount of residual penetration resulting from different fractions of the complete jet. The points at which the curves meet signify that the number of particles producing the superimposed curves are also the same number of particles producing the total residual penetration at that point. For example, PENJET's predictions at 11.5 CD standoff imply that predicting penetration using 40 particles would be just as accurate as using 51 particles. Anything less than 40 particles at that standoff would, however, produce inaccurate results. For the 24 CD standoff at which the round was actually fired, less than 20 digitized particles would be required for an accurate description of penetration. At short standoff, however, even 51 particles do not completely describe the penetration capabilities of Round 2331. Estimates are that below 8 CD standoff, particles beyond particle 51 were also capable of adding to penetration. For better quality rounds, though, as illustrated by Round 2937 in Figure 23, more particles are required even at longer standoffs to accurately describe penetration. For example, more than 50 digitized particles would be required for standoffs out to 25 CD for this particular round.

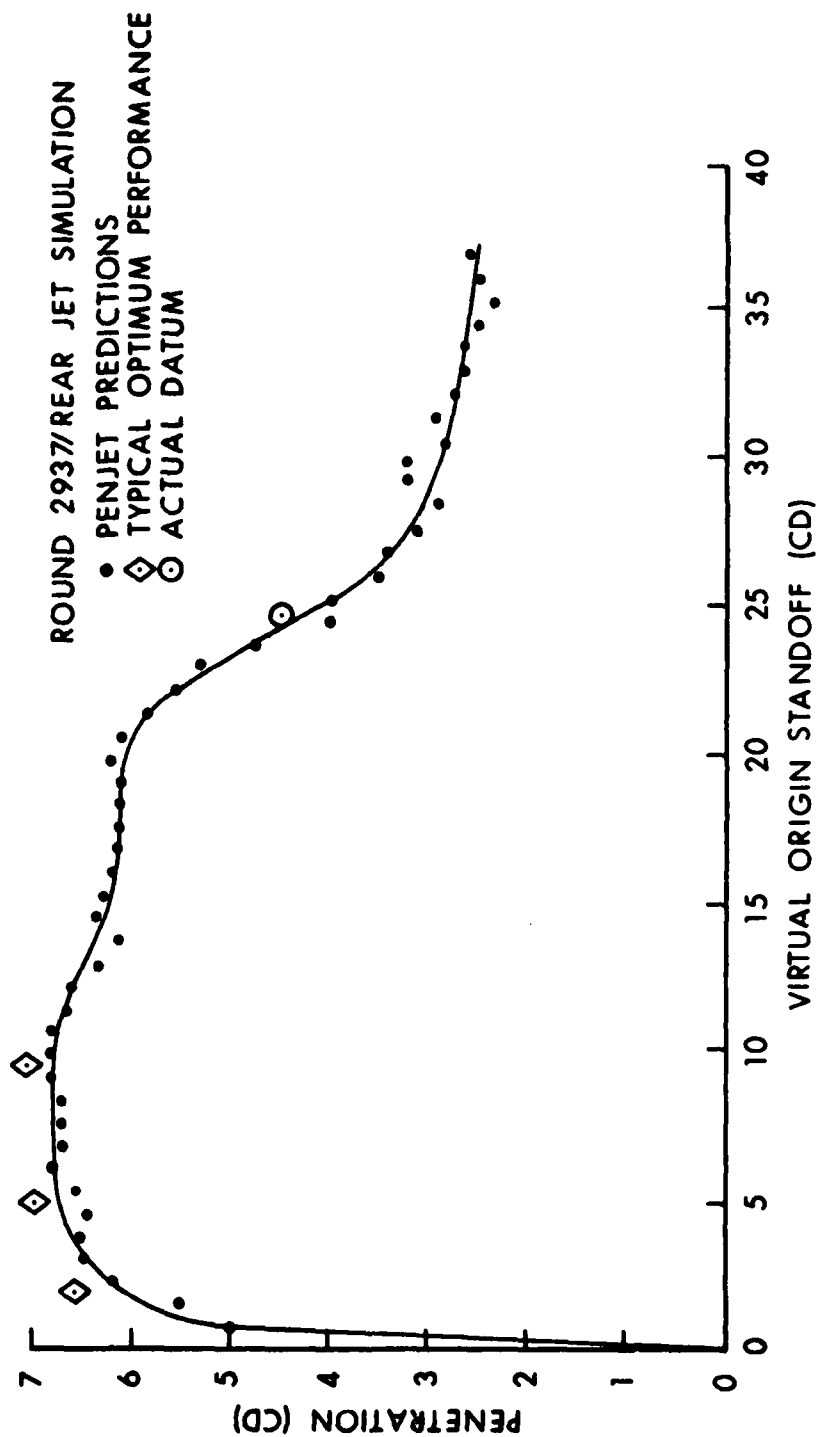


Figure 21. Analysis of Round 2937 with the Included Jet Rear Simulation
 Note: Optimum Performance Data Transferred from Figure 20.

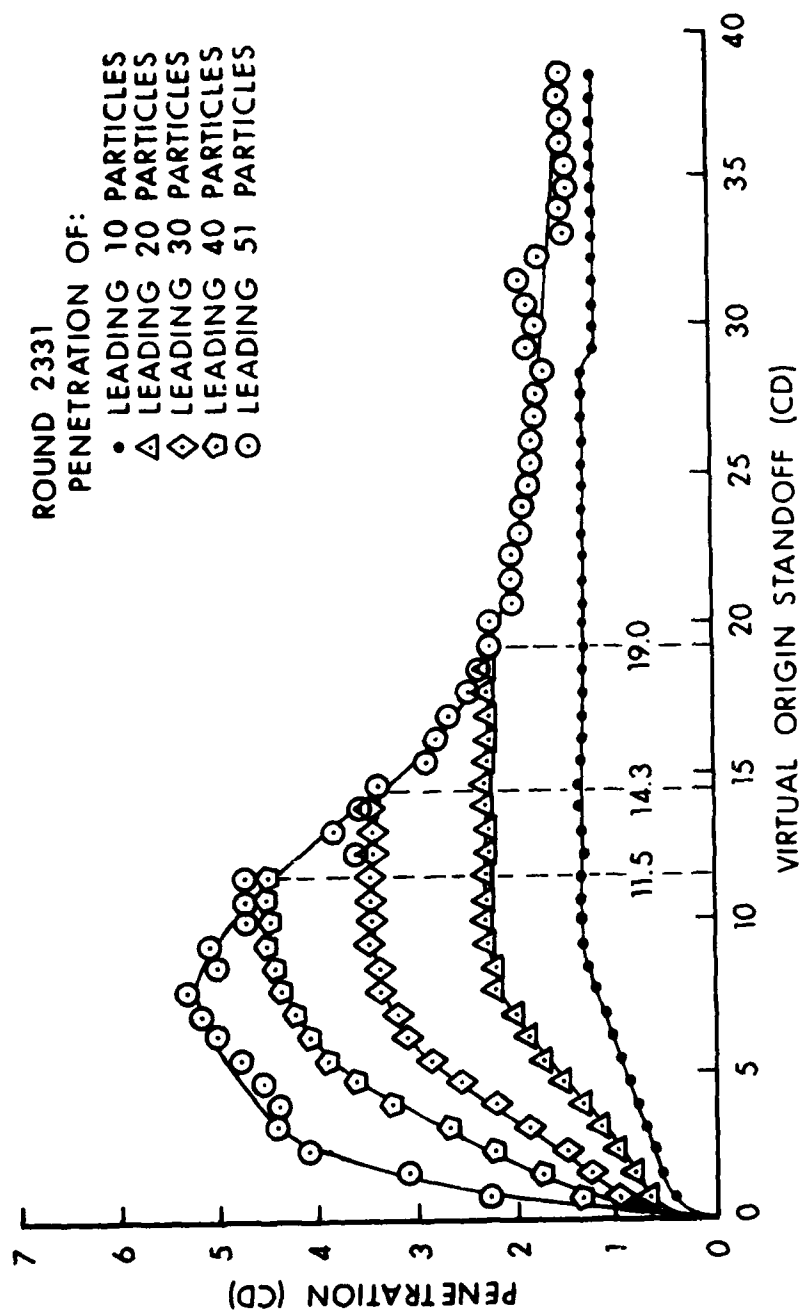


Figure 22. Penetration Capability of Distinct Jet Portions of Round 2331

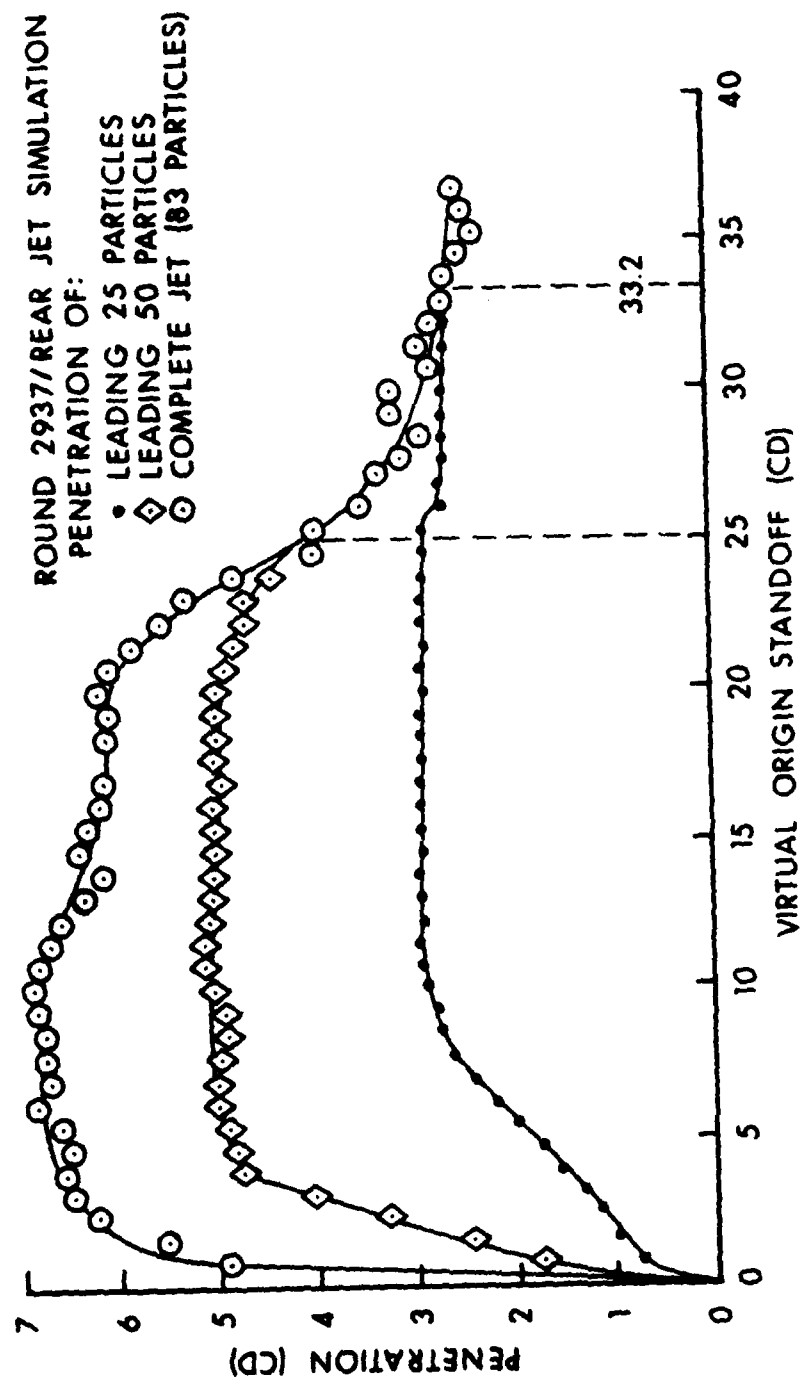


Figure 23. Penetration Capability of Distinct Jet Portions of Round 2937

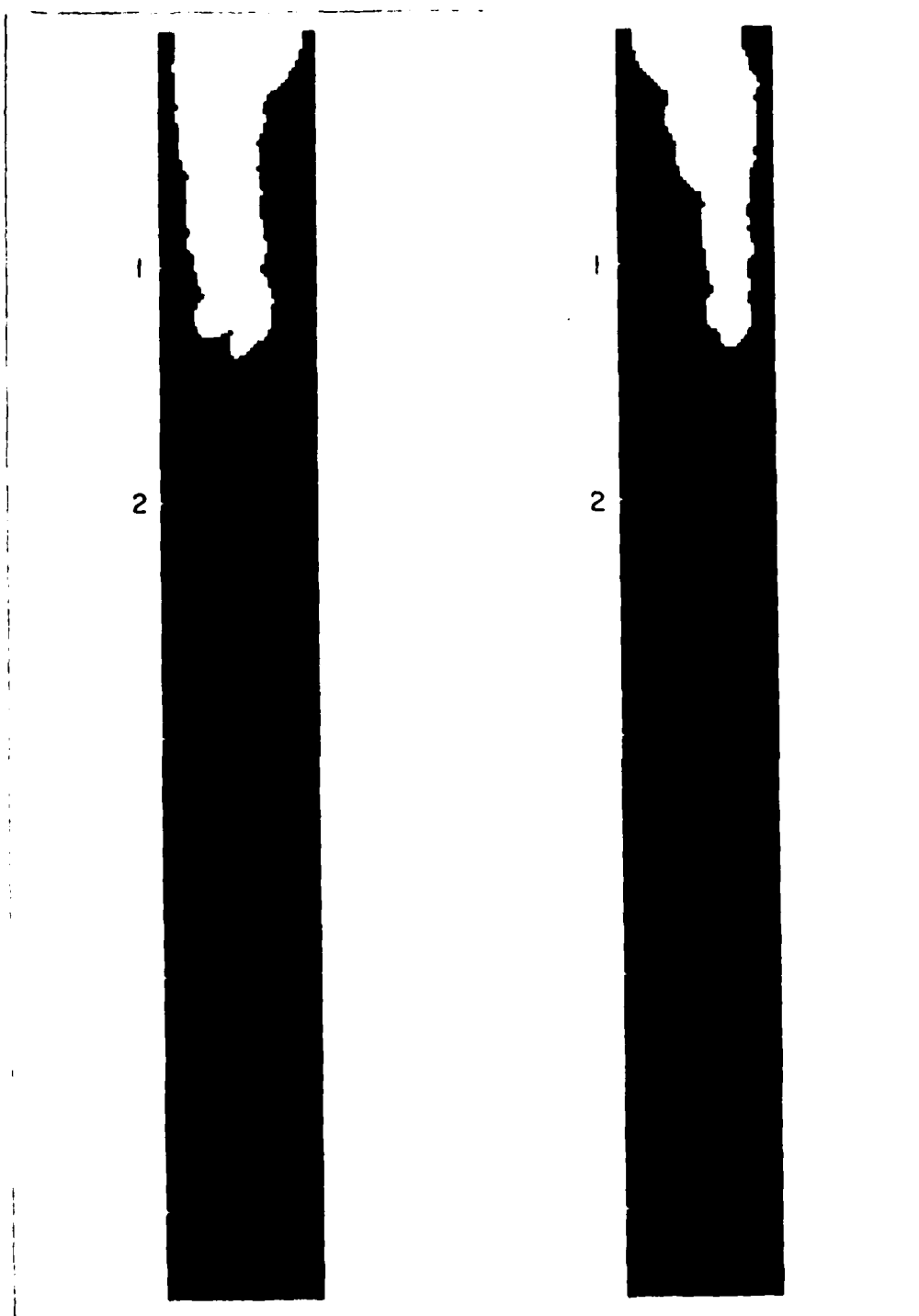


Figure 24. PENJET Predictions of Hole Profile for Round 2334 at 23 CD Standoff. Penetration in Cone Diameters.

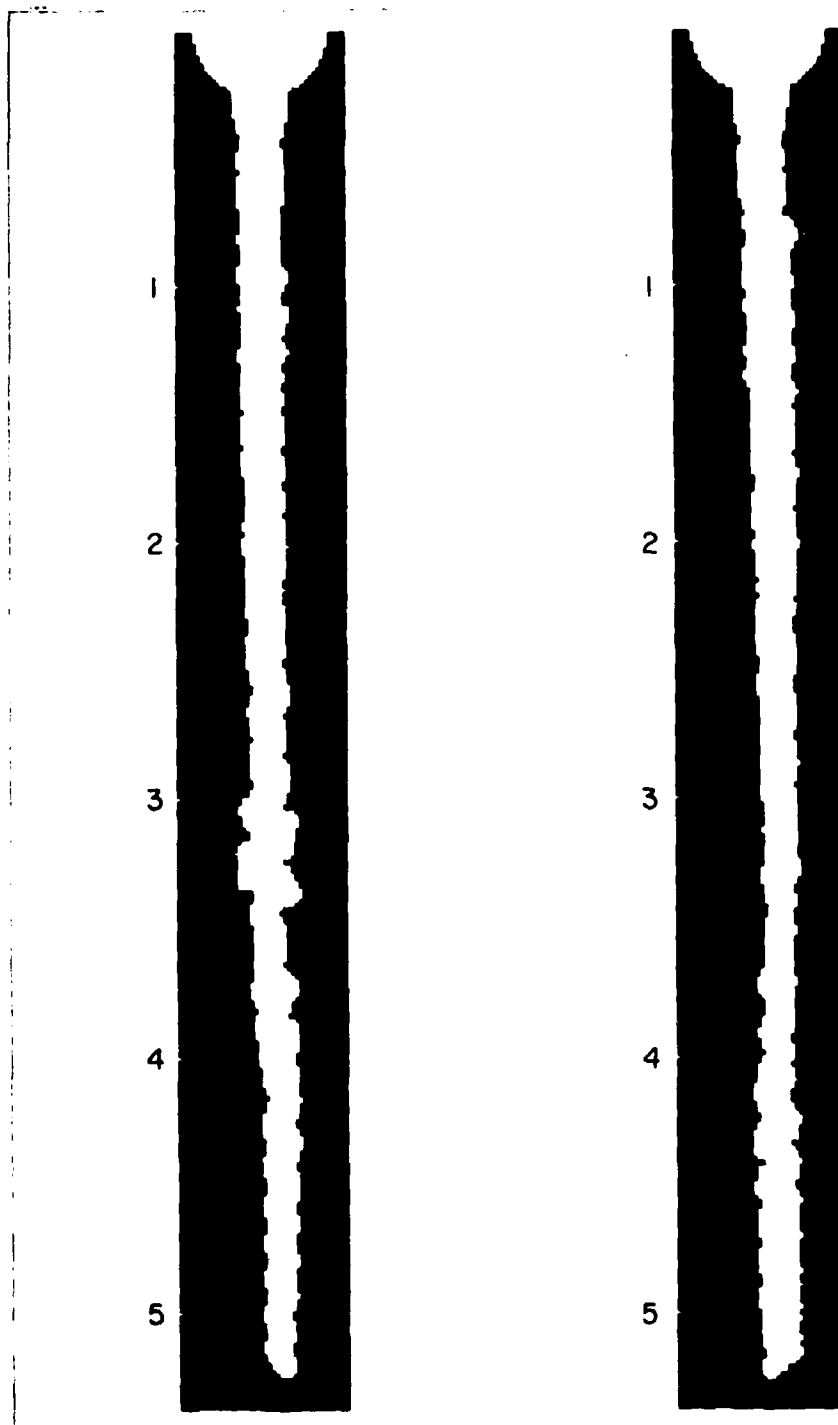


Figure 25. PENJET Predictions of Hole Profile for Round 2937 at 23 CD Standoff. Penetration in Cone Diameters.

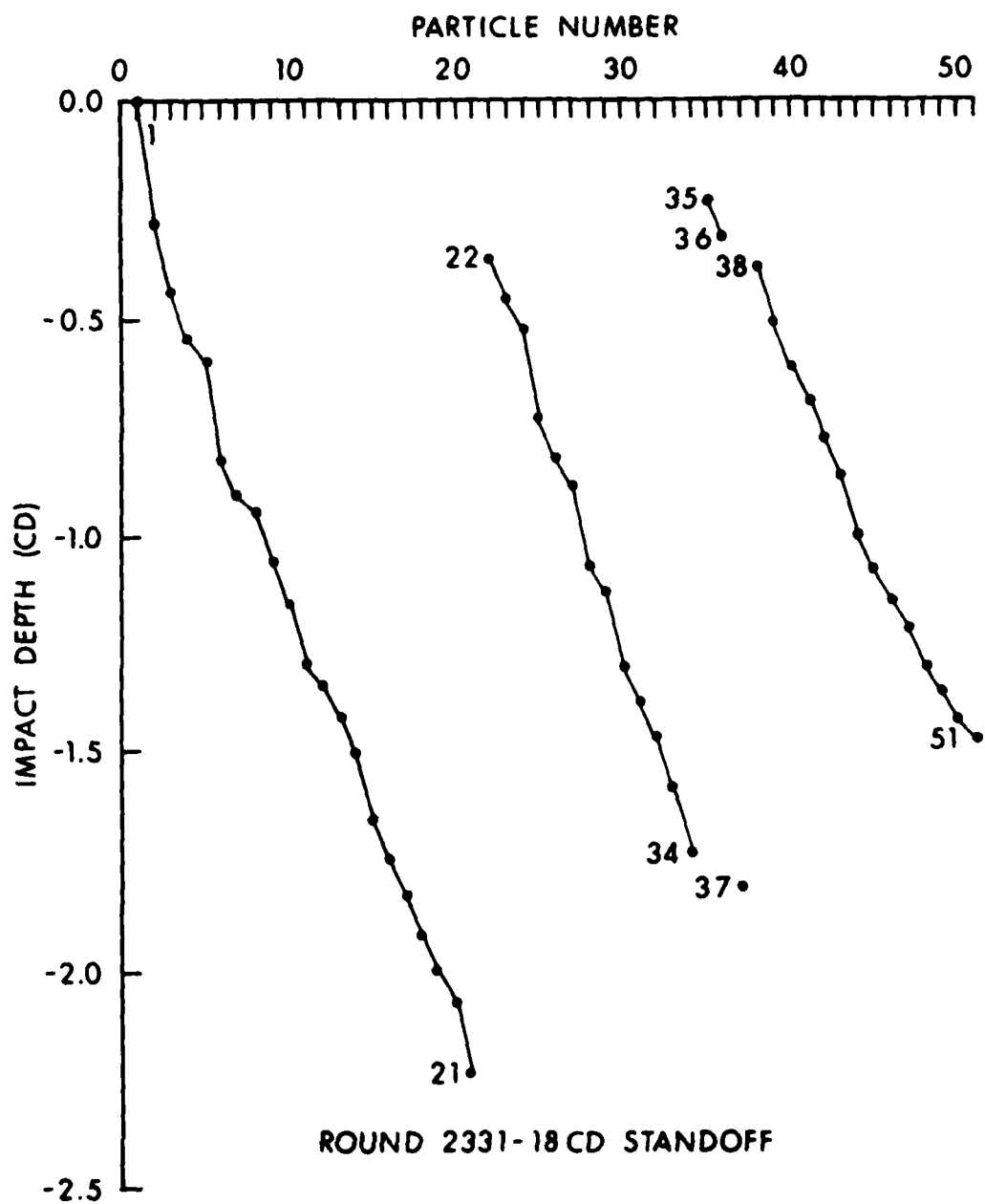


Figure 26. Code Predictions of Rear-Jet Ineffectiveness Due to Side Wall Impacts of Drifting Particles

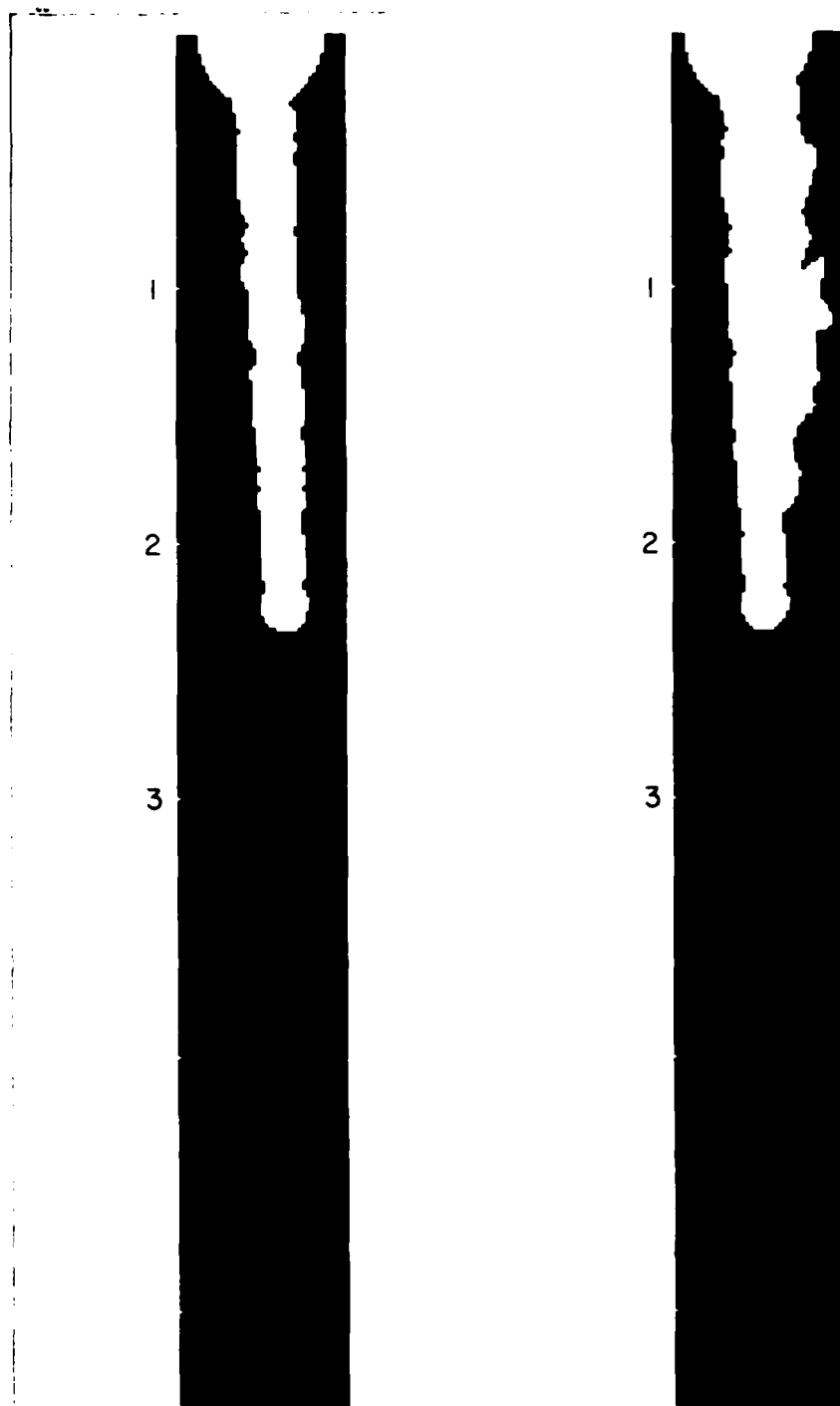


Figure 27. Orthogonal Views of Projected Hole Profile of Round 2331 at 18 CD Standoff as Predicted by PENJET. Penetration in Cone Diameters.

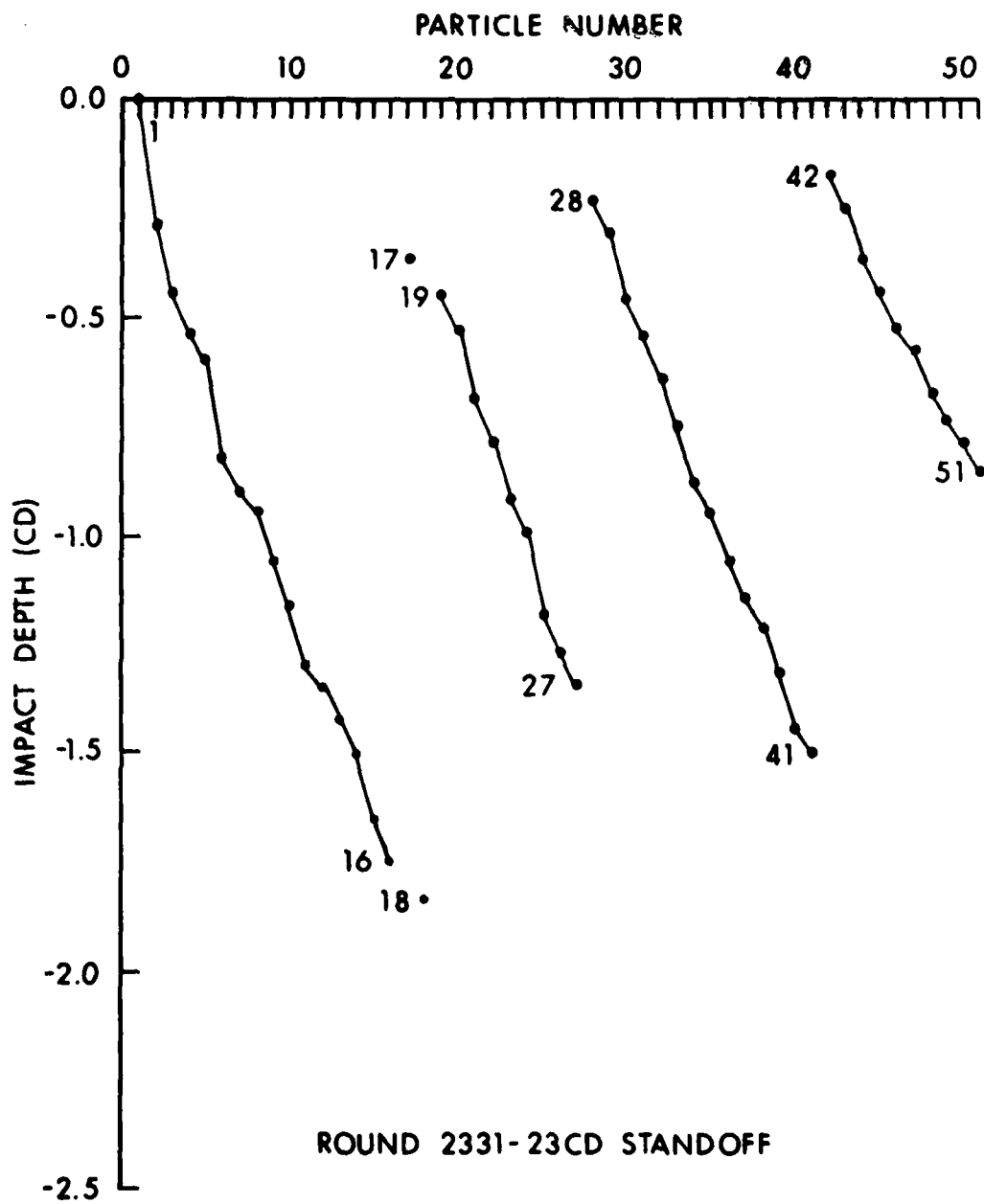


Figure 28. Increase of Side-Wall Impacts Predicted at Longer Standoffs

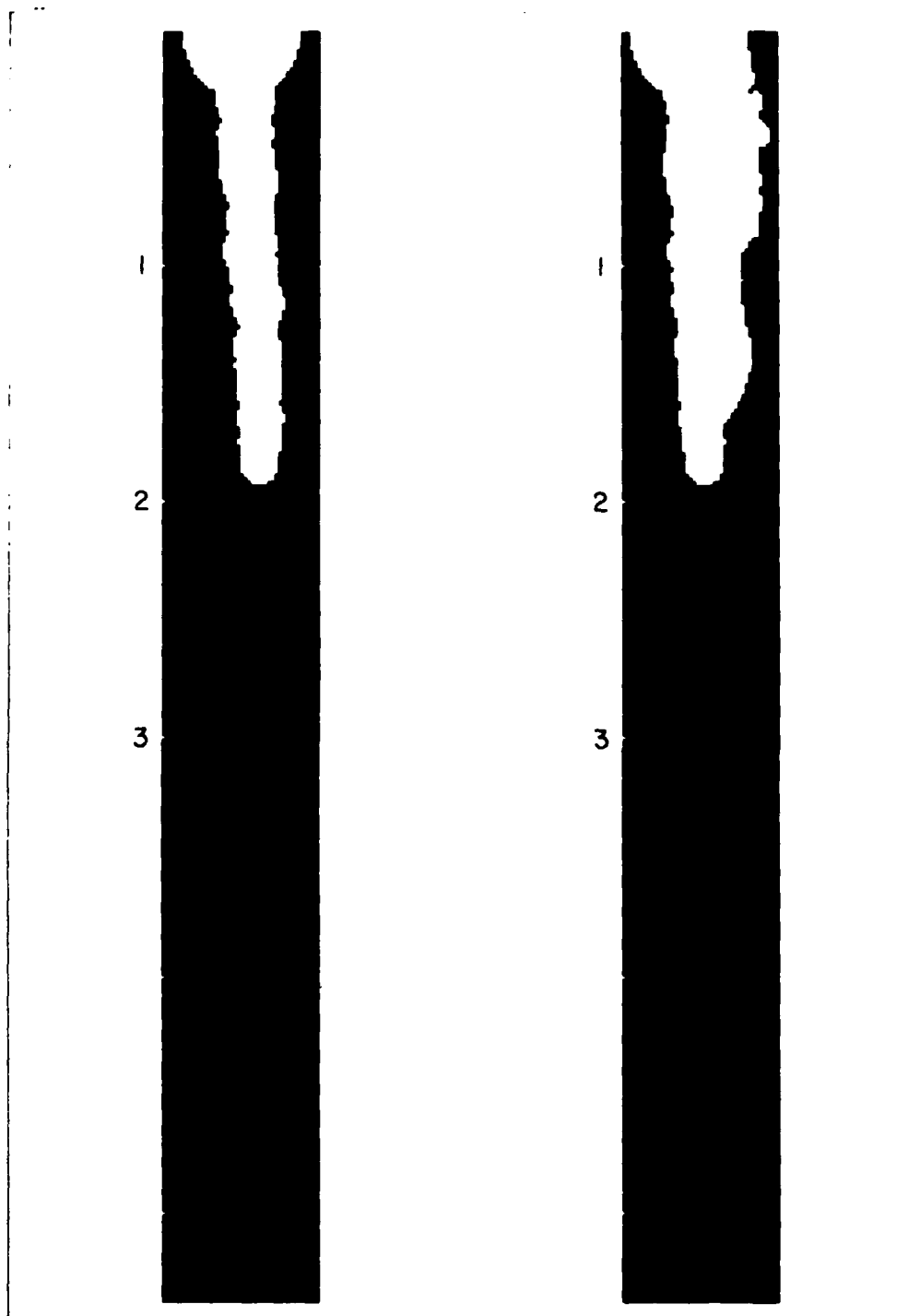


Figure 29. Orthogonal Views of Projected Hole Profile of Round 2331 at 23 CD Standoff as Predicted by PENJET. Penetration in Cone Diameters.

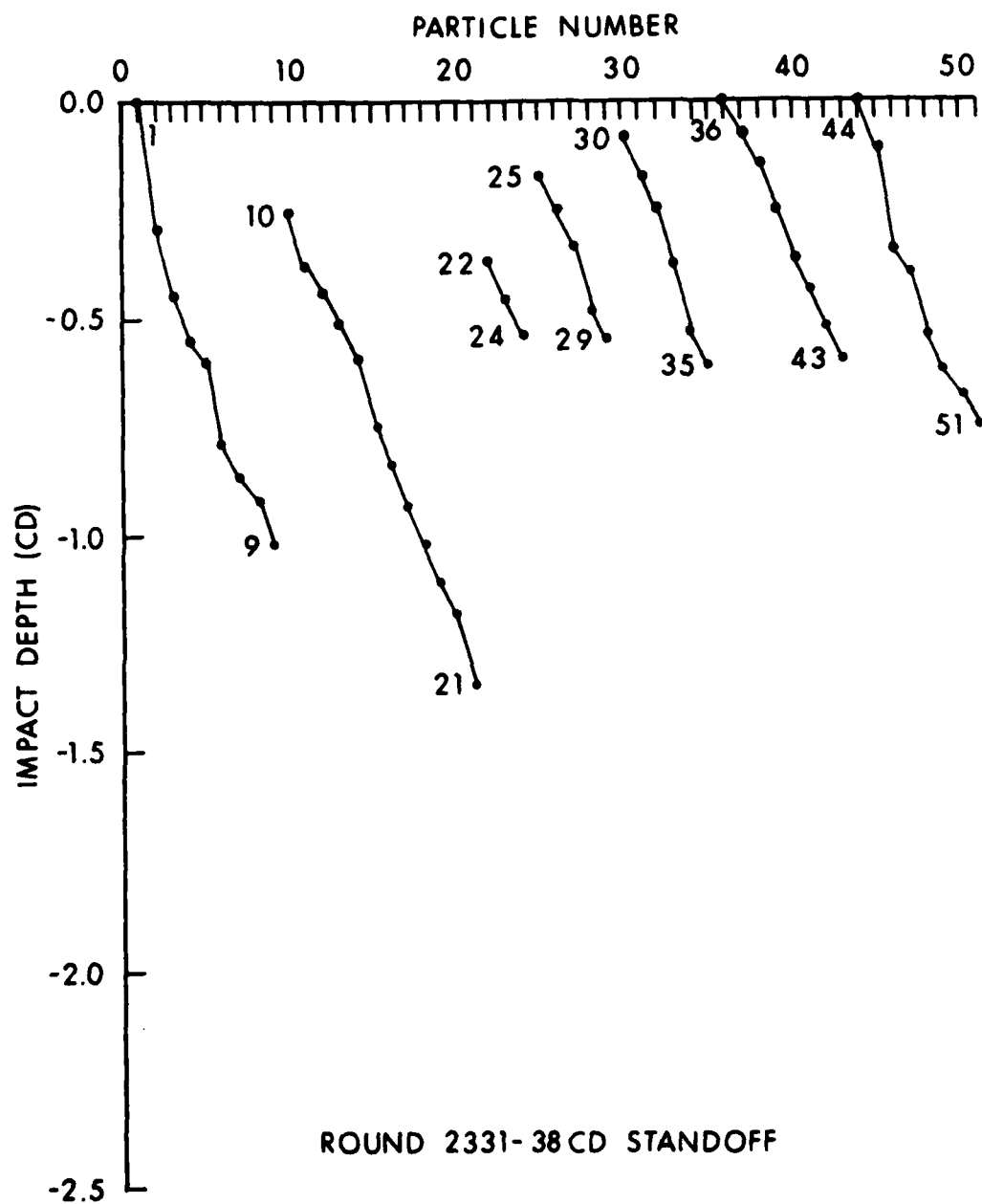


Figure 30. The Effect of Side Wall Impacts is Not Only to Decrease Penetration, But to Widen Crater.

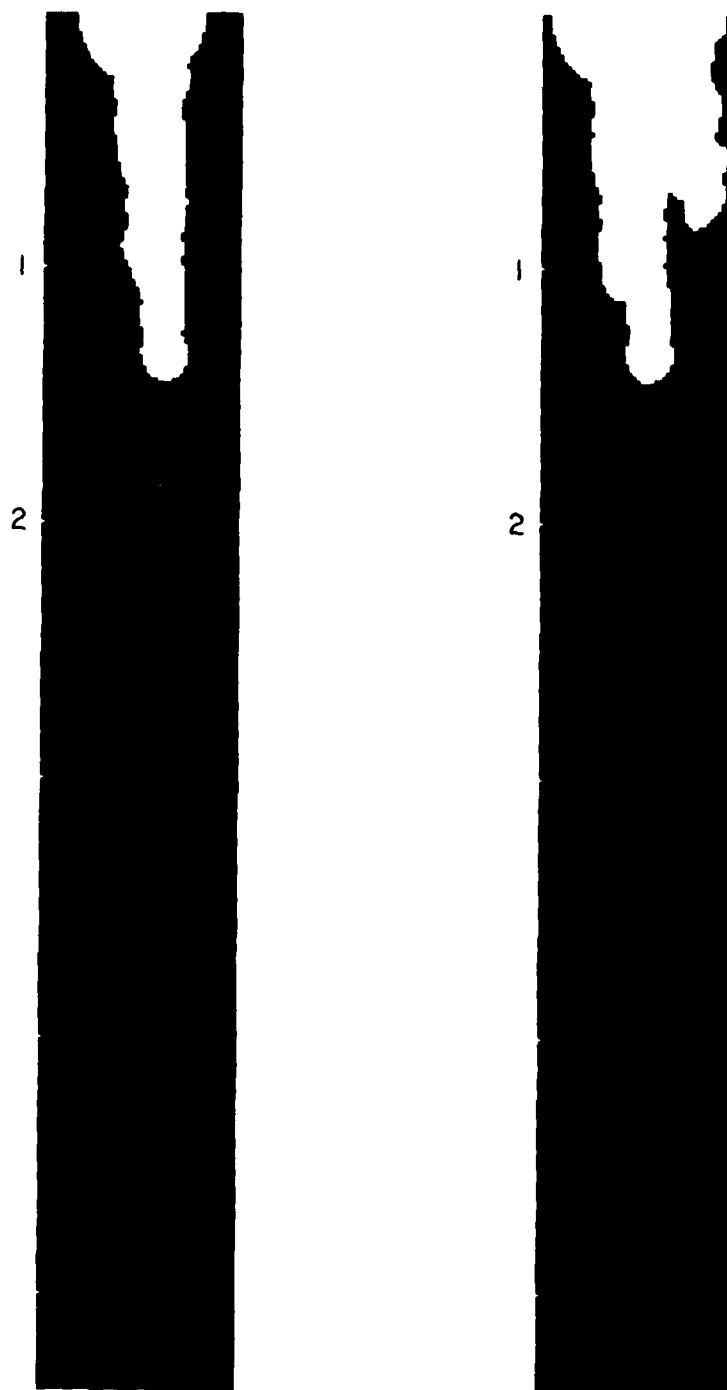


Figure 31. Orthogonal Views of Projected Hole Profile of Round 2331 at 38 CD Standoff as Predicted by PENJET. Penetration in Cone Diameters.

Besides showing a semblance of agreement between code and experiment, Figures 16-19 show that drift alone can easily account for rapid decreases in penetration as is often experienced by rounds at extended standoffs. The process by which penetration is decreased is one of sidewall impacts. This process was briefly mentioned before as a phenomenon which sometimes renders a particle ineffectual. The majority of sidewall impacts, though, are not so oblique as to disrupt penetration; most sidewall impacts go into creating wider hole profiles.

As an example of the devastating effects of drift velocity on jet penetration, a comparison of PENJET's hole profile predictions is shown for two separate rounds in Figures 24 and 25. The plots represent orthogonal views of the target block as described by PENJET. Both sets of calculations were performed at the same standoff (23 CD). The stark contrast of jet performance is noted not only in terms of depth of penetration, but also in terms of crater width. Sidewall impacts do not occur until much later in the penetration process for Round 2937, thus allowing more of the jet to contribute to the overall penetration.

Figure 26 is PENJET's prediction as to where in the hole profile each particle of Round 2331 strikes the target block at 18 CD standoff. Note that particle number is in order of formation. (Jet tip is particle 1.) It is seen that once a particle is blocked from reaching the bottom of the hole, it will usually begin the widening process somewhere up the profile. Subsequent particles try to work their way back down again. This process continues until no jet is left with which to penetrate the target. A plot of the hole profile as generated by PENJET is shown in Figure 27 as two orthogonal views of the target block. Graphs and plots for the same round are shown again at longer standoffs in Figures 28-31. It is seen that penetration is progressively lower because sidewall impacts occur earlier in the penetration process. Consequently, crater diameter progressively increases at longer standoffs. Thus, for Round 2331 at a standoff of 23 CD the total penetration capability of the jet as predicted by PENJET is contained within the first 18 particles.

The profile at 38 CD sharply illustrates the effects of drift velocity on jet penetration. The spread of the particles was so great as to cause the beginnings of a dual crater within 1 CD of the target surface. Thus, the ineffectiveness of the jet rear is seen. Any nondigitized portion of Round 2331 would probably produce its penetration in the vicinity of the particle 51 crater, far from the point of maximum penetration.

Thus, the decreases in penetration as predicted by PENJET are due mainly to the presence of drift velocities among the particles. Drift causes scattering of the particles and eventually, at some critical standoff, elimination of the particles from the crater deepening process via the mechanism of sidewall impacts.

VII. CONCLUSIONS

The ability to compute drift velocities for the particles of a shaped-charge jet is thought to be significant. In its own right, quantification of drift velocity aids in understanding the nature of the nonideal jet. A degree of imperfection may also be placed upon the jet in hopes of classifying a jet with regard to its quality.

The concept of how a jet penetrates is better understood when drift velocity is considered in the penetration picture. The process of sidewall impacts as a hinderance to penetration can be verified by inspection of actual target blocks. This process is reminiscent of the jet cutoff velocity which varies with standoff as proposed by DiPersio, Simon, and Merendino,³ in that the velocity of the particle striking deepest in the hole profile can be considered a cutoff velocity with respect to jet penetration. Implications here are that a jet could approach ideal penetration had drift velocity not been present. Larger craters would also make jets more insensitive to drift. Weaker targets allow for a wider crater for a given jet particle. Thus, the reduced probability of sidewall impact would allow greater penetration.

Compared to armor, a target of mild steel hydrodynamically should not affect the penetration of a given jet. It will affect penetration, however, when drift velocity is present in the jet. DiPersio and Simon⁸ verified this experimentally by firing jets into steel targets of varying hardness. Increasing crater size could alternately be accomplished by means of increasing the jet's energy. If dynamic impact did not induce additional drift in the jet, projecting the round into the target, as is the case under field conditions, should produce greater penetration than the static detonation of the same round. This cannot be verified experimentally, however, due mainly to the complexity of dynamic impact. Besides possessing possible yaw and pitch, the dynamic round is moving with respect to the target during the jet formation process. This relative motion can manifest itself in the form of drift. Dynamically detonated rounds do, however, create wider craters than their static counterparts as expected.

It is suggested that future investigations attempt to analyze drift from the viewpoint of warhead causes. The first step in understanding penetration in terms of drift has hopefully been set forth in this report. The ability to analyze drift as a function of the warhead would be the link between the imperfect warhead and nonideal penetration. Eventually, the possibility exists where warhead tolerances may be used to place limits upon the drift velocity distributions for a round so that penetration might be expressible as a function of charge tolerances.

⁸R. DiPersio, J. Simon, "The Effect of Target Hardness on the Penetration Capability of Shaped-Charge Jets," BRL Report No. 1408, July 1968 (AD 838991).

ACKNOWLEDGEMENTS

The author wishes to acknowledge the following members and emeriti of the Warhead Mechanics Branch of the U.S. Army Ballistic Research Laboratory for their encouraging support of this analysis: R. Jameson, M. Lampson, and J. Simon for technical consultations; V. Kucher for assistance rendered regarding computer related procedures; M. Scarborough and B. Simmons for having radiographs digitized and ready for use as computer data.

REFERENCE LIST

1. J. Blische, B. Simmons, "A Method for Reducing Data from Radiographs of Shaped-Charge Jets," BRL Report ARBRL-TR-02330, Jun 81 (AD A102770).
2. K. N. Kreyenhagen, J. E. Ferguson, R. R. Randall, J. P. Joyce, "Special Explosive Projectors," Proc. 6th Symposium Hypervelocity Impact, Vol. I (Held Cleveland, Ohio, April 30-May 2, 1963).
3. R. DiPersio, J. Simon, A. Merendino, "Penetration of Shaped-Charge Jets Into Metallic Targets," Ballistic Research Laboratory Report No. 1296, September 1965 (AD 476717).
4. E. Fitzgerald, Particle Waves and Deformation in Crystalline Solids, Interscience Publishers, New York, 1966.
5. J. Kineke, Jr. "An Experimental Study of Crater Formation in Metallic Targets," Proc. 4th Symposium Hypervelocity Impact, Vol. I (Held Eglin AFB, Florida, April 26-28, 1960).
6. G. Weihrauch, "The Behavior of Copper Pins Upon Impacting Various Materials with Velocities Between 50 m/s and 1650 m/s," (Doctoral Thesis) University of Karlsruhe, 1971.
7. J. Simon, R. DiPersio, "Experimental Verification of Standoff Effects on Shaped-Charge Jet Cutoff in Solid Targets," Ballistic Research Laboratory Memorandum No. 1976, May 1969 (AD 854396).
8. R. DiPersio, J. Simon, "The Effect of Target Hardness on the Penetration Capability of Shaped-Charge Jets," BRL Report No. 1408, July 1968 (AD 838991).

APPENDIX A

PROGRAM LISTING
AND
VARIABLE NAME DESCRIPTIONS

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The drift velocity and PENJET programs were written as direct additions to Blische's and Simmons' radiograph-data-reduction program (Reference 1) introduced in BRL Report 2330. Therefore, a listing of PENJET must be accompanied by the program which generates data for use within PENJET. The right-hand edge of the listing reveals where the statements in question originated. "GETREADY" are nonexecutable statements required for BRL Report 2330, PENJET, or both. "BRL 2330" is the bulk of Blische's and Simmons' program with minor revisions for PENJET applicability. "JETBREAK" is a revision for jet breakup calculations. "DRIFT" is used to calculate the drift distribution for a round while "PENJET" is the jet penetration program. The program is written in FORTRAN IV. An explanation of the program variables is inserted after the program listing.

It should be noted that even though Blische's and Simmons' program is versatile enough to run using data from only two X-ray flash exposures, PENJET is geared to the standard 3 flash setup and will require modifications to run otherwise.

Explanation for the program variable names found in the program segment identified as "BRL 2330" can be found in BRL Report 2330 (Reference 1). Should one wish to identify variables within the program segment "GETREADY," the other segments should be consulted, as all "GETREADY" variables are contained within other segments. All other significant variables are listed under their appropriate program segment.

PROGRAM MAIN(INPUT,OUTPUT,TAPES=INPUT,TAPE6=OUTPUT)	GETREADY	2
DIMENSION S1(100),S2(100),S3(100),VOL(100),XL(100),DIA(100),	GETREADY	3
1V1(100),V2(100),V3(100),VEL(100),XMAS3(100),XKE(100),SUMKE(100),	GETREADY	4
2ELOC(100),XXL(100),BURL(100),IX(6),IY(6),IZ(6),ZY(6),	GETREADY	5
3BREAK(100),SUMMAS(100),B(10000),DA(3),DB(3),XMA(3)	GETREADY	6
DIMENSION AA(2,3),CC(2),RR(300),AF(300),STG(2),TT(2)	GETREADY	7
DIMENSION DS(3),THETA(100),VETRAV(100),T1(100),T2(100),T3(100),	GETREADY	8
1 REF(3)	GETREADY	9
DIMENSION AVST(100),BETA(100)	GETREADY	10
DIMENSION PH1(100),PH2R(100),VTABS(100)	GETREADY	11
DIMENSION EPS(6,100),WT(6)	GETREADY	12
DIMENSION VX(100),VY(100)	GETREADY	13
DIMENSION R(2)	GETREADY	14
DIMENSION VT8(100)	GETREADY	15
DIMENSION T1(100),X(100),Y(100),Z(100),U(100),RADIUS(100),PEN(100),	GETREADY	16
1 HVOL(100)	GETREADY	17
DIMENSION M'C(100)	GETREADY	18
DIMENSION ZSUB(100),PENT(100),PENMAX(100)	GETREADY	19
DIMENSION V1(100,3),VM(100),VA(100,3),VAM(100),VP(100,3),VPM(100),	GETREADY	20
1 RGRAD(100,3),RGM(100),GRAD(100,3),AR(100,3),PR(100,3)	GETREADY	21
DIMENSION RL(100),RVEL(100)	GETREADY	22
DIMENSION OL(100),KVAL(100),TAU(100)	GETREADY	23
DIMENSION TP(100),RD(100)	GETREADY	24
REAL KD12,KD23,KD13	GETREADY	25
REAL INCR	GETREADY	26
DIMENSION SA(10),SB(10),SC(10),SD(10),SE(10),SF(10),SL1(10),	GETREADY	27
1SL2(10),SL3(10),ST(10),L1(100),L2(100),L3(100)	GETREADY	28
DIMENSION P(100),TOTP(100),SUMDTA(100),SUNLEN(100),XDTA(100),DELV(GETREADY	29
100),SDELV(100),SUMDEL(100)	GETREADY	30
DIMENSION XVOL(3,100)	GETREADY	31
DIMENSION AZ(100,100)	GETREADY	32
DATA SA(1),SA(2),SA(3)/10HVELOCITY (,10HMM/MICROSE,3HC)/	GETREADY	33
DATA SB(1),SB(2),SB(3)/10HCUMULATIVE,10H MASS (GRA,4HMS)/	GETREADY	34
DATA SC(1),SC(2),SC(3),SC(4)/10HPOSITION A,10HLONG JET L,10HENGTH	GETREADY	35
1(HM),1H)/	GETREADY	36
DATA SD(1),SD(2),SD(3)/10HCUMULATIVE,10H K.E. (JOU,5HLES)/	GETREADY	37
DATA SE(1),SE(2),SE(3),SE(4)/10HDISTANCE F,10HROM CHARGE,10H BASE	GETREADY	38
1(HM),1H)/	GETREADY	39
DATA SF(1),SF(2),SF(3)/10HBREAK-UP T,10HME (MICRO,5HSEC)/	GETREADY	40
DATA SL1(1)/8HFLASH 1>/	GETREADY	41
DATA SL2(1)/8HFLASH 2>/	GETREADY	42
DATA SL3(1)/8HFLASH 3>/	GETREADY	43
DATA MPA,MPAR,MPART,MCO,MCON,MCONT/4HPART,4HCUL,4HATED,4H CON,4HT	GETREADY	44
1INU,4HOU5 /	GETREADY	45
1 FORMAT(3I5,F3.2,6F10.5)	BRL2330	2
2 FORMAT(7F10.5)	BRL2330	3
6 FORMAT(8F10.5)	BRL2330	4
21 FORMAT(2E5)	BRL2330	5
30 FORMAT(12I5,2X,I1,2X,3I5)	BRL2330	6
READ(5,21) ICASES,LCL	BRL2330	7
IF(EOF(5)) 23,23	BRL2330	8
23 DO 500 IJ=1,ICASES	BRL2330	9
22 READ(5,11NROUND,MPART,MFLASH,RHC,(XMA(1),I=1,3),FLASH1,FLASH2,	BRL2330	10
1FLASH3	BRL2330	11
IF(EOF(5)) 24,24	BRL2330	12
24 ENCODE(21,20,ST(1)) NROUND	BRL2330	13
20 FORMAT(10HROUND NUMB,4HER ,I5,2H >)	BRL2330	14
READ(5,2) FOCUS1,FOCUS2,FOCUS3,F1A,F2A,F3A,F4A	BRL2330	15
IF(EOF(5)) 25,25	BRL2330	16
25 READ(5,6) F1B,F2B,F3B,F4B,F1C,F2C,F3C,F4C	BRL2330	17
IF(EOF(5)) 26,26	BRL2330	18
26 RD=RHO	BRL2330	19
IF(LCL.NE.1) GO TO 29	BRL2330	20

RHO=0.	BRL2330	21
C	BRL2330	22
ROUND INFORMATION CARD NEXT	BRL2330	23
29 PRINT 5	BRL2330	24
PRINT 3, NROUND, RHO, (XMAG(I), N=1, 3), FOCUS1, FOCUS2, FOCUS3, FLASH1, FL	BRL2330	25
1ASH2, FLASH3	BRL2330	26
3 FORMAT(//, //, 20X, 'ROUND NUMBER ', 15, //, 20X, 'LINER DENSITY (GM/CC)-	BRL2330	27
1 ', F6.1, //, 20X, 'MAGNIFICATION FACTOR-', 3F9.3, //, 20X, 'DISTANCE FROM L	BRL2330	28
ZINER BASE TO FOCAL POINT (MM)', //, 25X, 'FLASH 1- ', F6.1, //, 25X, 'FLASH	BRL2330	29
32- ', F6.1, //, 25X, 'FLASH 3- ', F6.1, //, 20X, 'DELAY TIMES (MICROSEC)', //,	BRL2330	30
425X, 'FLASH 1- ', F6.1, //, 25X, 'FLASH 2- ', F6.1, //, 25X, 'FLASH 3- ', F6.1	BRL2330	31
5)	BRL2330	32
RHO=RO	BRL2330	33
PRINT 5	BRL2330	34
5 FORMAT(1H1)	BRL2330	35
PRINT 4	BRL2330	36
4 FORMAT(21X, 'PARTICLE AVG. VELOCITY TOTAL JET BREAK-UP', //, 2	BRL2330	37
12X, 'NUMBER', 4X, ' (MM/MICROSEC)', 3X, 'LENGTH (MM)', 3X, ' (PICROSEC)', //)	BRL2330	38
P1=3.14159	BRL2330	39
DO 150 I=1, NFLASH	BRL2330	40
DO 149 J=1, NPART	BRL2330	41
READ(5, 30) (IX(K), IY(K), K=1, 6), L, IFILM, IPDUND, IPART	BRL2330	42
IF (EOF(5)) 32, 32	BRL2330	43
32 DO 35 M=1, 6	BRL2330	44
ZX(M)=FLOAT(IX(M))/15.5906	BRL2330	45
35 ZY(M)=FLOAT(IY(M))/15.5906	BRL2330	46
AZ(1, 8)=((ZY(2)+ZY(5))/2.)*XMAG(1)	BRL2330	47
IF (ZX(8).EQ. ZX(2)) GO TO 40	BRL2330	48
R1=.5*SORT((ZX(1)-ZX(6))*2+(ZY(1)-ZY(6))*2)*XMAG(1)	BRL2330	49
R2=.5*SORT((ZX(2)-ZX(5))*2+(ZY(2)-ZY(5))*2)*XMAG(1)	BRL2330	50
R3=.5*SORT((ZX(3)-ZX(4))*2+(ZY(3)-ZY(4))*2)*XMAG(1)	BRL2330	51
P1X=(ZX(1)+ZX(6))/2.	BRL2330	52
P1Y=(ZY(1)+ZY(6))/2.	BRL2330	53
P2X=(ZX(3)+ZX(4))/2.	BRL2330	54
P2Y=(ZY(3)+ZY(4))/2.	BRL2330	55
P3X=(ZX(2)+ZX(5))/2.	BRL2330	56
P3Y=(ZY(2)+ZY(5))/2.	BRL2330	57
XW1=SORT((P1X-P3X)*2+(P1Y-P3Y)*2)*XMAG(1)	BRL2330	58
XW2=SORT((P3X-P2X)*2+(P3Y-P2Y)*2)*XMAG(1)	BRL2330	59
P1Z=P1X	BRL2330	60
P1X=P2X	BRL2330	61
GOTO 45	BRL2330	62
40 P1X=ZX(1)	BRL2330	63
45 IF (L.EQ. 2) GOTO 55	BRL2330	64
IF (L.EQ. 3) GOTO 65	BRL2330	65
IF (IFILM.EQ. 2) GO TO 47	BRL2330	66
IF (IFILM.EQ. 3) GO TO 48	BRL2330	67
IF (IFILM.EQ. 4) GO TO 49	BRL2330	68
S1(J)=FOCUS1-(F1A-P1X)*XMAG(1)	BRL2330	69
GO TO 15	BRL2330	70
47 S1(J)=FOCUS1-(F2A-P1X)*XMAG(1)	BRL2330	71
GO TO 15	BRL2330	72
48 S1(J)=FOCUS1+(F3A+P1X)*XMAG(1)	BRL2330	73
GO TO 15	BRL2330	74
49 S1(J)=FOCUS1+(F4A+P1X)*XMAG(1)	BRL2330	75
15 IF (ZX(1).EQ. ZX(2)) GO TO 50	BRL2330	76
L1(J)=0	BRL2330	77
GOTO 75	BRL2330	78
50 L1(J)=1	BRL2330	79
GOTO 80	BRL2330	80
55 IF (IFILM.EQ. 2) GO TO 57	BRL2330	81
IF (IFILM.EQ. 3) GO TO 58	BRL2330	82
IF (IFILM.EQ. 4) GO TO 59	BRL2330	83
S2(J)=FOCUS2-(F1B-P1X)*XMAG(1)	BRL2330	84

GO TO 16	BRL2330	85
57 S2(J)=FOCUS2-(F2B-P1X)*XMAG(I)	BRL2330	86
GO TO 16	BRL2330	87
58 S2(J)=FOCUS2-(F3B-P1X)*XMAG(I)	BRL2330	88
GO TO 16	BRL2330	89
59 S2(J)=FOCUS2-(F4B-P1X)*XMAG(I)	BRL2330	90
16 V1(J)=(S2(J)-S1(J))/(FLASH2-FLASH1)	BRL2330	91
IF(ZX(1).EQ.ZX(2)) GO TO 60	BRL2330	92
L2(J)=0	BRL2330	93
GOTO 75	BRL2330	94
60 L2(J)=1	BRL2330	95
GOTO 80	BRL2330	96
65 IF(NFLASH.EQ.2) GOTO 80	BRL2330	97
IF(FILM.EQ.2) GO TO 67	BRL2330	98
IF(FILM.EQ.3) GO TO 68	BRL2330	99
IF(FILM.EQ.4) GO TO 69	BRL2330	100
S3(J)=FOCUS3-(F1C-P1X)*XMAG(I)	BRL2330	101
GO TO 17	BRL2330	102
67 S3(J)=FOCUS3-(F2C-P1X)*XMAG(I)	BRL2330	103
GO TO 17	BRL2330	104
68 S3(J)=FOCUS3-(F3C-P1X)*XMAG(I)	BRL2330	105
GO TO 17	BRL2330	106
69 S3(J)=FOCUS3-(F4C-P1X)*XMAG(I)	BRL2330	107
17 V2(J)=(S3(J)-S2(J))/(FLASH3-FLASH2)	BRL2330	108
V3(J)=(S3(J)-S1(J))/(FLASH3-FLASH1)	BRL2330	109
IF(ZX(1).EQ.ZX(2)) GO TO 70	BRL2330	110
L3(J)=0	BRL2330	111
GOTO 75	BRL2330	112
70 L3(J)=1	BRL2330	113
GOTO 80	BRL2330	114
75 VOL(J)= VOL(J)+(PI*XM1/3.*(R1**2+R1*R2+R2**2)+PI*XM2/3.*(R2**2+R2*	BRL2330	115
R3+R3**2))	BRL2330	116
P1X=P1Z	BRL2330	117
XVOL(I,J)= PI*XM1/3.*(R1**2+R1*R2+R2**2)+PI*XM2/3.*(R2**2+R2*R3+	BRL2330	118
R3**2)	BRL2330	119
XVOL(I,J)= XVOL(I,J)*.001/RHO	BRL2330	120
XL(I)=XL(I)+SORT((P1X-P2X)**2+(P1Y-P2Y)**2)*XMAG(I)	BRL2330	121
DIA(J)=DIA(J)+2.*R2	BRL2330	122
80 CONTINUE	BRL2330	123
IF(I.LT.NFLASH) GOTO 149	BRL2330	124
IF(NFLASH.EQ.2) GOTO 85	BRL2330	125
IF(L1(J).EQ.1.AND.L2(J).EQ.1.AND.L3(J).EQ.1) GO TO 149	BRL2330	126
IF(L1(J).EQ.0.AND.L2(J).EQ.0.AND.L3(J).EQ.0) GO TO 90	BRL2330	127
IF(L2(J).EQ.1.AND.L3(J).EQ.1) GO TO 95	BRL2330	128
IF(L1(J).EQ.1.AND.L3(J).EQ.1) GO TO 95	BRL2330	129
IF(L1(J).EQ.1.AND.L2(J).EQ.1) GO TO 95	BRL2330	130
IF(L2(J).EQ.0.AND.L3(J).EQ.0) GO TO 100	BRL2330	131
IF(L1(J).EQ.0.AND.L3(J).EQ.0) GO TO 100	BRL2330	132
IF(L1(J).EQ.0.AND.L2(J).EQ.0) GO TO 100	BRL2330	133
85 IF(L1(J).EQ.1.AND.L2(J).EQ.1) GO TO 149	BRL2330	134
IF(L1(J).EQ.0.AND.L2(J).EQ.0) GO TO 100	BRL2330	135
IF(L1(J).EQ.0.OR.L2(J).EQ.0) GO TO 95	BRL2330	136
90 XMASS(J)=VOL(J)*.001/3.*RHO	BRL2330	137
XL(U)=XL(J)/3.	BRL2330	138
DIA(J)=DIA(J)/3.	BRL2330	139
GOTO 149	BRL2330	140
95 XMASS(U)=VOL(J)*.001/RHO	BRL2330	141
GOTO 149	BRL2330	142
100 XMASS(J)=VOL(J)*.001/2.*RHO	BRL2330	143
XL(J)=XL(I)/2.	BRL2330	144
DIA(J)=DIA(J)/2.	BRL2330	145
149 CONTINUE	BRL2330	146
150 CONTINUE	BRL2330	147
	JETBREAK	1

C

C	BROK IS THE NUMBER OF JET PARTICLES THAT HAVE PARTICULATED AS OF	JETBREAK	2
C	EXPOSURE TIMES 1, 2, AND 3 RESPECTIVELY. TOLAY IS THE DELAY TIME	JETBREAK	3
C	BETWEEN DETONATION AND JET FORMATION (MICROSEC).	JETBREAK	4
C		JETBREAK	5
	READ(5,*) BROK1,BROK2,BROK3,TOLAY	JETBREAK	6
	PZ=BROK1+BROK2+BROK3	JETBREAK	7
	TZ=FLASH1+FLASH2+FLASH3	JETBREAK	8
	PBYTZ=BROK1*FLASH1+BROK2*FLASH2+BROK3*FLASH3	JETBREAK	9
	PBYTZ=BROK1*2+BROK2*2+BROK3*2	JETBREAK	10
	Z1=(PBYTZ-(PZ*TZ/3.))/(PBYTZ-(PZ*2/3.))	JETBREAK	11
	PBAR=PZ/3.	JETBREAK	12
	TAVG=TZ/3.	JETBREAK	13
	DO 170 N=1,NPART	BRL2330	148
	IF(NFLASH.EQ.2)GOTO 195	BRL2330	149
	VEL(N)=(V1(N)+V2(N)+V3(N))/3.*.1	BRL2330	150
	GOTO 160	BRL2330	151
155	VEL(N)=V1(N)*.1	BRL2330	152
160	SUML(N)=XL(N)+SUML(N-1)	BRL2330	153
	SUMMAS(N)=XMASS(N)+SUMMAS(N-1)	BRL2330	154
	ELOD(N)=XL(N)/DIA(N)	BRL2330	155
	VEL(N)=VEL(N)*10.	BRL2330	156
	XKE(N)=.5*XMASS(N)*.001*(VEL(N)*1000.)*2	BRL2330	157
	SUMKE(N)=XKE(N)+SUMKE(N-1)	BRL2330	158
	P(N)=VEL(N)*XMASS(N)	BRL2330	159
	TOTP(N)=P(N)+TOTP(N-1)	BRL2330	160
	SUMDIA(N)=SUMDIA(N-1)+DIA(N)	BRL2330	161
	BREAK(N)=TAVG+Z1*(N-PBAR)-TOLAY	JETBREAK	14
	GOTO 170	BRL2330	162
165	BREAK(1)=0.0	BRL2330	163
170	PRINT 171, N,VEL(N),SUML(N),BREAK(N)	BRL2330	164
171	FORMAT(23X,T2,10X,F6.3,9X,F6.2,7X,F6.1)	BRL2330	165
	PRINT 92	BRL2330	166
	92 FORMAT(1H1,20X,'PARTICLE',4X,'VELOCITY1',6X,'VELOCITY2',6X,	BRL2330	167
	1'VELOCITY3',/,22X,'NUMBER',3X,'(MM/MICROSEC)',2X,'(MM/MICROSEC)',	BRL2330	168
	22X,'(MM/MICROSEC)',/)	BRL2330	169
	DO 172 J=1,NPART	BRL2330	170
172	PRINT 173, J,V1(J),V2(J),V3(J)	BRL2330	171
173	FORMAT(23X,T2,9X,F6.3,7X,F6.3,9X,F6.3)	BRL2330	172
	PRINT 175	BRL2330	173
175	FORMAT(1H1,20X,'PARTICLE LENGTH DIA. L/D MASS TOTAL J	BRL2330	174
	1ET',/,22X,'NUMBER',3X,'(MM)',4X,'(MM)',9X,'(GRAMS)',2X,'MASS(GRAMS	BRL2330	175
	2)',/)	BRL2330	176
	DO 176 I=1,NPART	BRL2330	177
176	PRINT 177, I,XL(I),DIA(I),ELOD(I),XMASS(I),SUMMAS(I)	BRL2330	178
177	FORMAT(23X,T2,8X,F4.1,4X,F4.1,2X,F4.1,4X,F9.2,5X,F6.2)	BRL2330	179
	PRINT 603	BRL2330	180
603	FORMAT(1H1,20X,'PARTICLE',4X,'MASS1',9X,'MASS2',9X,'MASS3',/,22X,	BRL2330	181
	1'NUMFR',4X,'(GRAMS)',7X,'(GRAMS)',7X,'(GRAMS)',/)	BRL2330	182
	DO 600 J=1,NPART	BRL2330	183
	I=1	BRL2330	184
600	WRITE(6,602) (J,XVOL(I,J),XVOL(I+1,J),XVOL(I+2,J))	BRL2330	185
602	FORMAT(23X,T2,6X,F8.4,6X,F8.4,6X,F8.4)	BRL2330	186
	PRINT 180	BRL2330	187
180	FORMAT(1H1,20X,'PARTICLE K.E. TOTAL JET DISTANCE FROM CH	BRL2330	188
	1ARGE BASE',/,22X,'NUMBER',4X,'(JOULES)',3X,'KE(JOULES)',2X,'FLASH	BRL2330	189
	21',2X,'FLASH 2',2X,'FLASH 3',/)	BRL2330	190
	DO 181 I=1,NPART	BRL2330	191
181	PRINT 182, I,XKE(I),SUMKE(I),S1(I),S2(I),S3(I)	BRL2330	192
182	FORMAT(23X,T2,7X,F8.0,4X,F8.0,4X,F9.0,4X,F9.0)	BRL2330	193
	PRINT 183	BRL2330	194
183	FORMAT(1H1,20X,'PARTICLE',4X,'MOMENTUM',4X,'TOTAL JET',/,22X,'NUMB	BRL2330	195
	1ER',4X,'(KG-M/SEC)',3X,'MOMENTUM',/)	BRL2330	196
	DO 185 I=1,NPART	BRL2330	197
185	PRINT 184, I,P(I),TOTP(I)	BRL2330	198

184	FORMAT(23X,I2,9X,F6.2,7X,F6.2)	BRL2330	199
	DO 192 N=2,NPART	BRL2330	200
	SUMLEN(N)=SUMLEN(N-1)+XL(N)	BRL2330	201
	XDIA(N)=XDIA(N-1)+DIA(N)	BRL2330	202
	IF(N.EQ.2) GO TO 191	BRL2330	203
	DELV(N)=VEL(N-1)-VEL(N)	BRL2330	204
	GO TO 192	BRL2330	205
191	DELV(N)=VEL(1)-VEL(N)	BRL2330	206
192	SDELV(N)=SDELV(N-1)+DELV(N)	BRL2330	207
	DO 195 J=3,NPART	BRL2330	208
195	SUNDEL(J)=DELV(J)+SUNDEL(J-1)	BRL2330	209
	AVL1=SUMLEN(NPART)/FLOAT(NPART)	BRL2330	210
	AVL2=SUMLEN(NPART)/FLOAT(NPART-1)	BRL2330	211
	AVD1=SUMDIA(NPART)/FLOAT(NPART)	BRL2330	212
	AVD2=XDIA(NPART)/FLOAT(NPART-1)	BRL2330	213
	ADELV1=SDELV(NPART)/FLOAT(NPART-1)	BRL2330	214
	ADELV2=SUNDEL(NPART)/FLOAT(NPART-2)	BRL2330	215
	PRINT 200, AVL1, AVL2, AVD1, AVD2, ADELV1, ADELV2	BRL2330	216
200	FORMAT(1H1,7777,47X,'WITH JET TIP W/D TIP',/,20X,'AVERAGE PART	BRL2330	217
	ICLE LENGTH',7X,F6.2,7X,F6.2,/,20X,'AVERAGE PARTICLE DIAMETER',6X,	BRL2330	218
	2F5.2,8X,F5.2,/,20X,'AVERAGE CHANGE IN VELOCITY',6X,F4.2,9X,F4.2)	BRL2330	219
	PRINT 905	BRL2330	220
	NRA=2	BRL2330	221
	NN=1	BRL2330	222
	TC=0	BRL2330	223
	CALL POLYLS(VEL,S1,NPART,AA,NRA,NN,CC,RR,AF,ERMS,SIG,TT,DET,IC)	BRL2330	224
	WRITE(6,41) CC(1)	BRL2330	225
	DS(1)=CC(1)	BRL2330	226
41	FORMAT(20X,' VIRTUAL ORIGIN FOR FLASH 1-',F12.6)	BRL2330	227
	CALL POLYLS(VEL,S2,NPART,AA,NRA,NN,CC,RR,AF,ERMS,SIG,TT,DET,IC)	BRL2330	228
	WRITE(6,42) CC(1)	BRL2330	229
	DS(2)=CC(1)	BRL2330	230
42	FORMAT(20X,' VIRTUAL ORIGIN FOR FLASH 2-',F12.6)	BRL2330	231
	IF(L.LT.3) GO TO 515	BRL2330	232
	CALL POLYLS(VEL,S3,NPART,AA,NRA,NN,CC,RR,AF,ERMS,SIG,TT,DET,IC)	BRL2330	233
	WRITE(6,43) CC(1)	BRL2330	234
	DS(3)=CC(1)	BRL2330	235
43	FORMAT(20X,' VIRTUAL ORIGIN FOR FLASH 3-',F12.6)	BRL2330	236
C		DRIFT	1
C	DRIFT CALCULATIONS	DRIFT	2
C		DRIFT	3
C		DRIFT	4
C	EACH REF REFERS TO THE DISTANCE BETWEEN THE REFERENCE PARTICLE	DRIFT	5
C	AND THE FIDUCIAL FROM WHERE ALL OTHER MEASUREMENTS ARE TAKEN.	DRIFT	6
C	REF IS TAKEN IN MM AND SHOULD BE THE ACTUAL DISTANCE READ OFF OF	DRIFT	7
C	THE FILM. IT WILL THEN BE MULTIPLIED BY THE MAG FACTOR(WNECH	DRIFT	8
C	SHOULD HAVE BEEN ENTERED AS DATA IN BRL2330) NOTE, IF XMAG IS	DRIFT	9
C	LESS THAN 1. , THEN THE FILM PORTRAYS AN IMAGE LARGER THAN REALITY	DRIFT	10
C		DRIFT	11
	READ(5,*) REF(1),REF(2),REF(3)	DRIFT	12
	REF(1)=REF(1)*XMAG(1)	DRIFT	13
	REF(2)=REF(2)*XMAG(2)	DRIFT	14
	REF(3)=REF(3)*XMAG(3)	DRIFT	15
	DO 401 J=1,3	DRIFT	16
	DO 400 I=1,NPART	DRIFT	17
400	AZ(I,J)=AZ(I,J)-REF(J)	DRIFT	18
401	CONTINUE	DRIFT	19
	PRINT 33	DRIFT	20
33	FORMAT(1H1,21X,'PARTICLE',8X,'DEVIANCE FROM PATH (MM)',/,23X,	DRIFT	21
	1'NUMBER',6X,'FLASH 1 FLASH 2 FLASH 3',/)	DRIFT	22
	DO 800 JJJ=1,NPART	DRIFT	23
	WRITE (6,36) JJJ,AZ(1,JJJ),AZ(2,JJJ),AZ(3,JJJ)	DRIFT	24
36	FORMAT(24X,I2,8X,F8.3,3X,F8.3,3X,F8.3)	DRIFT	25
800	CONTINUE	DRIFT	26

DO 491 I=1,NPART	DRIFT	27
T1(I)=(S1(I)-DS(1))/VEL(I)	DRIFT	28
T2(I)=(S2(I)-DS(2))/VEL(I)	DRIFT	29
T3(I)=(S3(I)-DS(3))/VEL(I)	DRIFT	30
KD12=AZ(1,I)*T2(I)/(AZ(2,I)*T1(I))	DRIFT	31
KD23=AZ(2,I)*T3(I)/(AZ(3,I)*T2(I))	DRIFT	32
KD13=AZ(1,I)*T3(I)/(AZ(3,I)*T1(I))	DRIFT	33
DO 405 J=1,3	DRIFT	34
IF(AZ(IU,I)+REF(J)) 405,402,405	DRIFT	35
402 THETA(I)=0.0	DRIFT	36
VETRAV(I)=0.0	DRIFT	37
GO TO 491	DRIFT	38
405 CONTINUE	DRIFT	39
C SOLVING FOR THETA USING FIRST TWO GOVERNING EQUATIONS	DRIFT	40
C	DRIFT	41
TH12=ATAN(2.0*KD12/SQRT(2.0)-1.0)	DRIFT	42
IF(TH12.LT.0.0)TH12=TH12+PI	DRIFT	43
C SOLVING FOR THETA USING LAST TWO GOVERNING EQUATIONS	DRIFT	44
C	DRIFT	45
TH23=ATAN(1.0-2.0/(SQRT(2.0)*KD23))	DRIFT	46
IF(TH23.LT.0.0)TH23=TH23+PI	DRIFT	47
450 IF(ABS(TH12-TH23).LT.PI/2.0)GO TO 449	DRIFT	48
IF(TH23.LT.TH12)GO TO 448	DRIFT	49
TH12=TH12+PI	DRIFT	50
GO TO 449	DRIFT	51
448 TH12=TH12-PI	DRIFT	52
C SOLVING FOR THETA USING FIRST AND LAST GOVERNING EQUATIONS	DRIFT	53
C	DRIFT	54
449 TH13=ATAN((KD13-1.0)/(KD13+1.0))	DRIFT	55
IF(TH13.LT.0.0)TH13=TH13+PI	DRIFT	56
470 IF(ABS(TH23-TH13).LT.PI/2.0)GO TO 469	DRIFT	57
IF(TH23.LT.TH13)GO TO 468	DRIFT	58
TH13=TH13+PI	DRIFT	59
GO TO 469	DRIFT	60
468 TH13=TH13-PI	DRIFT	61
C FINDING AN AVERAGE VALUE FOR THETA AND PLACING IT IN THE PROPER	DRIFT	62
C QUADRANT	DRIFT	63
469 THV=(TH12+TH23+TH13)/3.0	DRIFT	64
IF((TH23-PI/2.0)*AZ(2,I)) 471,473,472	DRIFT	65
471 THETA(I)=THV	DRIFT	66
GO TO 480	DRIFT	67
472 THETA(I)=THV+PI	DRIFT	68
TH12=TH12+PI	DRIFT	69
TH23=TH23+PI	DRIFT	70
TH13=TH13+PI	DRIFT	71
IF(THETA(I)-THV.GT.1.7)THV=THV+PI	DRIFT	72
GO TO 480	DRIFT	73
473 IF(TH13) 474,475,475	DRIFT	74
474 THETA(I)=THV+PI	DRIFT	75
TH12=TH12+PI	DRIFT	76
TH23=TH23+PI	DRIFT	77
TH13=TH13+PI	DRIFT	78
IF(THETA(I)-THV.GT.1.7)THV=THV+PI	DRIFT	79
GO TO 480	DRIFT	80
475 THETA(I)=THV	DRIFT	81
C CALCULATING RELATIVE TRANSVERSE VELOCITIES WHILE ASSIGNING	DRIFT	82
C WEIGHTED VALUES...(DEPENDENT ON THE PARTICLES' ORIENTATION	DRIFT	83
C WITH THE FLASH X-RAY)	DRIFT	84
	DRIFT	85
	DRIFT	86
	DRIFT	87
	DRIFT	88
	DRIFT	89
	DRIFT	90

C	480	VETR12=AZ(1,I)/(COS(TH12-PI/4.0)*T1(I))*1000.0	DRIFT	91
		VETR23=AZ(2,I)/(COS(TH23)*T2(I))*1000.0	DRIFT	92
		VETR13=AZ(3,I)/(COS(TH13+PI/4.0)*T3(I))*1000.0	DRIFT	93
		EPS(1,I)=AZ(1,I)*SIN(TH12-PI/4.0)*(TH12-TMV)/T1(I)*1000.0	DRIFT	94
	1	/(COS(TH12-PI/4.0))*2	DRIFT	95
		EPS(2,I)=AZ(2,I)*SIN(TH23)*(TH23-TMV)*1000.0/T2(I)	DRIFT	96
	1	/(COS(TH23))*2	DRIFT	97
		EPS(3,I)=AZ(3,I)*SIN(TH13+PI/4.0)*(TH13-TMV)/T3(I)*1000.0	DRIFT	98
	1	/(COS(TH13+PI/4.0))*2	DRIFT	99
		EPS(4,I)=AZ(1,I)*SIN(TH13-PI/4.0)*(TH13-TMV)/T1(I)*1000.0	DRIFT	100
	1	/(COS(TH13-PI/4.0))*2	DRIFT	101
		EPS(5,I)=AZ(2,I)*SIN(TH12)*(TH12-TMV)*1000.0/T2(I)	DRIFT	102
	1	/(COS(TH12))*2	DRIFT	103
		EPS(6,I)=AZ(3,I)*SIN(TH23+PI/4.0)*(TH23-TMV)/T3(I)*1000.0	DRIFT	104
	1	/(COS(TH23+PI/4.0))*2	DRIFT	105
		DO 481 J=1,6	DRIFT	106
	481	EPS(J,I)=ABS(EPS(J,I))	DRIFT	107
		EPSUM=0.0	DRIFT	108
		DO 482 J=1,6	DRIFT	109
	482	EPSUM=1.0/EPS(J,I)*EPSUM	DRIFT	110
		DO 483 J=1,6	DRIFT	111
	483	WT(J)=1.0/(EPS(J,I)*EPSUM)	DRIFT	112
		VETRAV(I)=WT(1)*VETR12+WT(2)*VETR23+WT(3)*VETR13+WT(4)*VETR12	DRIFT	113
	1	+WT(5)*VETR23+WT(6)*VETR13	DRIFT	114
C			DRIFT	115
C		CONVERT THETA TO DEGREES	DRIFT	116
C			DRIFT	117
	490	THETA(I)=180.0*THETA(I)/PI	DRIFT	118
	491	CONTINUE	DRIFT	119
		PRINT 492	DRIFT	120
	492	FORMAT(1H,21X,'PARTICLE',5X,'RELATIVE TRANSVERSE',5X,'RELATIVE ',	DRIFT	121
	1	'DEVIANC',/,23X,'NUMBER',5X,'VELOCITY (M/S)',9X,'ANGLE',	DRIFT	122
	2	' (DEGREES)',/)	DRIFT	123
		DO 495 K=1,NPART	DRIFT	124
		WRITE(6,494) K,VETRAV(K),THETA(K)	DRIFT	125
	494	FORMAT(24X,12,15X,F6.2,17X,F5.1)	DRIFT	126
	495	CONTINUE	DRIFT	127
C			DRIFT	128
C		TRANSLATE APPARENT VALUES INTO ABSOLUTE VALUES	DRIFT	129
C			DRIFT	130
		R(2)=0.0	DRIFT	131
C			DRIFT	132
C		GUESSING AT AN INITIAL DEVIANCE ANGLE FOR THE REFERENCE	DRIFT	133
C			DRIFT	134
		THMIN=THETA(I)	DRIFT	135
		THMAX=THETA(I)	DRIFT	136
		MINDEX=1	DRIFT	137
		MAXDEX=1	DRIFT	138
		DO 391 I=2,NPART	DRIFT	139
		IF(THETA(I).LT.THMIN.OR.THMIN-THETA(I).LT.-300.160 TO 392	DRIFT	140
		GO TO 393	DRIFT	141
	392	THMIN=THETA(I)	DRIFT	142
		MINDEX=I	DRIFT	143
		GO TO 391	DRIFT	144
	393	IF(THETA(I).GT.THMAX.OR.THMAX-THETA(I).GT.300.) GO TO 394	DRIFT	145
		GO TO 391	DRIFT	146
	394	THMAX=THETA(I)	DRIFT	147
		MAXDEX=I	DRIFT	148
	391	CONTINUE	DRIFT	149
		IF(MINDEX.GT.MAXDEX) GO TO 395	DRIFT	150
		INCR=.0034907	DRIFT	151
		THETA0=THETA(MAXDEX)-180.	DRIFT	152
		GO TO 396	DRIFT	153
			DRIFT	154

395	INCR=-.0034907	DRIFT	155
	THETA0=THETA(MINDEX)-180.	DRIFT	156
396	ORCS=0.0	DRIFT	157
	DO 899 I=1,NPART	DRIFT	158
	THETA(I)=THETA(I)+PI/180.0	DRIFT	159
	VX(I)=VETRAV(I)*COS(THETA(I))	DRIFT	160
899	VY(I)=VETRAV(I)*SIN(THETA(I))	DRIFT	161
	THOR=THETA0+PI/180.0	DRIFT	162
902	R(I)=R(I)	DRIFT	163
	ITOTL=IFIX(.122173/ABS(INCR))	DRIFT	164
	DO 1100 NOT=1,ITOTL	DRIFT	165
	THOR=THOR+INCR	DRIFT	166
903	VTO=0.0	DRIFT	167
C		DRIFT	168
C	FINDING THE MAGNITUDE OF REFERENCE TRANSVERSE VELOCITY THAT WILL	DRIFT	169
C	MAKE IT PROPORTIONAL TO THE AXIAL VELOCITY	DRIFT	170
C		DRIFT	171
	DO 756 I=1,NPART	DRIFT	172
756	VTABS(I)=VETRAV(I)	DRIFT	173
904	VTO=VTO+10.0	DRIFT	174
	INC=100	DRIFT	175
	GO TO 909	DRIFT	176
905	VTO=VTO+1.0	DRIFT	177
	INC=10	DRIFT	178
	GO TO 909	DRIFT	179
906	VTO=VTO+.1	DRIFT	180
	INC=1	DRIFT	181
909	VX0=VTO*COS(THOR)	DRIFT	182
	VY0=VTO*SIN(THOR)	DRIFT	183
	DO 907 I=1,NPART	DRIFT	184
907	VTABS(I)=SQRT((VX(I)-VX0)**2+(VY(I)-VY0)**2)	DRIFT	185
C		DRIFT	186
C	STATISTICAL CALCULATIONS FOR A LINEAR REGRESSION	DRIFT	187
C		DRIFT	188
802	TS=0.0	DRIFT	189
	AS=0.0	DRIFT	190
	TBYAS=0.0	DRIFT	191
	TBYTS=0.0	DRIFT	192
	ABYAS=0.0	DRIFT	193
	KLESS=0	DRIFT	194
	DO 900 I=1,NPART	DRIFT	195
	IF(ABS(VTABS(I)-VTO) .LE. 1.0) GO TO 700	DRIFT	196
	TS=VTABS(I)+TS	DRIFT	197
	AS=VEL(I)*AS	DRIFT	198
	TBYAS=VTABS(I)*VEL(I)+TBYAS	DRIFT	199
	TBYTS=VTABS(I)**2+TBYTS	DRIFT	200
	ABYAS=VEL(I)**2+ABYAS	DRIFT	201
	GO TO 900	DRIFT	202
700	KLESS=KLESS+1	DRIFT	203
900	CONTINUE	DRIFT	204
	CO1=(TBYAS-(TS*AS/(NPART-KLESS)))/(TBYTS-(TS**2/(NPART-KLESS)))	DRIFT	205
	CO2=(TBYAS-(TS*AS/(NPART-KLESS)))/(ABYAS-(AS**2/(NPART-KLESS)))	DRIFT	206
	CO=(CO1+1.0/CO2)/2.0	DRIFT	207
	TBAR=TS/(NPART-KLESS)	DRIFT	208
	ABAR=AS/(NPART-KLESS)	DRIFT	209
	XINT=TBAR-ABAR/CO	DRIFT	210
	IF(XINT.LT.0.0 .AND. INC.EQ.100) GO TO 904	DRIFT	211
	IF(XINT.LT.0.0 .AND. INC.EQ.10) GO TO 905	DRIFT	212
	IF(XINT.LT.0.0 .AND. INC.EQ.1) GO TO 906	DRIFT	213
	IF(INC=10) 912,911,910	DRIFT	214
910	VTO=VTO-11.0	DRIFT	215
	GO TO 905	DRIFT	216
911	VTO=VTO-1.1	DRIFT	217
	GO TO 906	DRIFT	218

912 R(2)=SQRT(C01+C02)	DRIFT	219
IF(R(2).GT.R(1)) GO TO 1101	DRIFT	220
GO TO 1100	DRIFT	221
1101 R(1)=R(2)	DRIFT	222
VTOB=VTO	DRIFT	223
THORB=THOR	DRIFT	224
DO 1102 IB=1,NPART	DRIFT	225
1102 VT0(IB)=VTABS(IB)	DRIFT	226
1100 CONTINUE	DRIFT	227
VTO=VTOB	DRIFT	228
THOR=THORB	DRIFT	229
DO 1103 IOI=1,NPART	DRIFT	230
1103 VTABS(IOI)=VT0(IOI)	DRIFT	231
700 DO 815 M=1,NPART	DRIFT	232
IF((VTABS(M)-VTO).LE. 1.0) GO TO 811	DRIFT	233
PHI(M)=THETA(M)-THOR	DRIFT	234
BETA(M)=ASIN(VECTRAV(M)*SIN(PHI(M))/VTABS(M))*180.0/PI	DRIFT	235
AVST(M)=VEL(M)/VTABS(M)*1000.0	DRIFT	236
GO TO 815	DRIFT	237
811 VTABS(M)=0.0	DRIFT	238
BETA(M)=0.0	DRIFT	239
AVST(M)=0.0	DRIFT	240
815 CONTINUE	DRIFT	241
PRINT 702	DRIFT	242
702 FORMAT(1H1,21X,'PARTICLE',5X,'ABSOLUTE TRANSVERSE',5X,'ABSOLUTE ',	DRIFT	243
1 'ANGULAR',5X,'AXIAL VELOCITY OVER',/,23X,'NUMBER',9X,'VEL',	DRIFT	244
2 'OCITY (M/S)',4X,'DEVANCE WRT REFERENCE TRANSVERSE VELOC',	DRIFT	245
3 'ITY',/)	DRIFT	246
DO 710 J=1,NPART	DRIFT	247
WRITE(6,707) J,VTABS(J),BETA(J),AVST(J)	DRIFT	248
707 FORMAT(24X,I2,16X,F6.2,19X,F6.1,17X,F6.2)	DRIFT	249
710 CONTINUE	DRIFT	250
PRENT 712,VTO	DRIFT	251
712 FORMAT(7,22X,'THE TRANSVERSE VELOCITY OF THE REFERENCE IS ',	DRIFT	252
1 F6.2,' M/S')	DRIFT	253
PRINT 713, R(2)	DRIFT	254
713 .DRMAT(7,22X,'THE CORRELATION COEFFICIENT FOR THIS REGRESSION',	DRIFT	255
1 ' IS ',F9.4)	DRIFT	256
C ***** P E N J E T *****	PENJET	1
C	PENJET	2
C	PENJET	3
C	PENJET	4
C ECONST IS THE ENERGY REQUIRED TO CREATE ONE CUBIC-MM OF HOLE	PENJET	5
C VOLUME. RHOJ IS JET DENSITY IN G/CU-MM. SIGJ IS JET STRENGTH IN	PENJET	6
C MEGAPASCALS.. RHOY AND SIGY ARE THE CORRESPONDING TARGET	PENJET	7
C VARIABLES	PENJET	8
C	PENJET	9
READ(5,300)ECONST,RHOJ,SIGJ,RHOY,SIGY	PENJET	10
300 FORMAT(9F10.4)	PENJET	11
C	PENJET	12
C ZMIN IS THE MINIMUM STANDOFF AT WHICH PENJET WILL EVALUATE THE	PENJET	13
C GIVEN DRIFT DISTRIBUTION. ZMAX IS THE MAXIMUM STANDOFF TO BE	PENJET	14
C ANALYZED. ZINC IS THE INCREMENT BETWEEN STANDOFF ANALYZATION.	PENJET	15
C ALL MEASUREMENTS ARE IN MM.	PENJET	16
C	PENJET	17
READ(5,397)ZMIN,ZMAX,ZINC	PENJET	18
397 FORMAT(3F10.1)	PENJET	19
DO 301 I=1,NPART	PENJET	20
BETA(I)=BETA(I)*3.14159/180.	PENJET	21
RADIUS(I)=0.0	PENJET	22
PEN(I)=0.0	PENJET	23
RL(I)=XL(I)	PENJET	24
RVEL(I)=VEL(I)	PENJET	25
301 CONTINUE	PENJET	26

NM1=NPART-1	PENJET	27
GAMMA=SQRT(RHOT/RHOJ)	PENJET	28
ZSUBO=ZMIN-ZINC	PENJET	29
KZSUM=0	PENJET	30
250 ZSUBO=ZSUBO+ZINC	PENJET	31
KZSUM=KZSUM+1	PENJET	32
PENMAX(KZSUM)=0.	PENJET	33
ZSUB(KZSUM)=ZSUBO	PENJET	34
C	PENJET	35
C THIS LOOP EVALUATES EACH DIGITIZED PARTICLE	PENJET	36
C	PENJET	37
291 DO 209 I=1,NPART	PENJET	38
NCC(I)=0	PENJET	39
NO=0	PENJET	40
260 T(I)=(ZSUBO-XL(I)-(DS(1)+DS(2)+DS(3))/3.-2.995)/VEL(I)	PENJET	41
DELT=3./VEL(I)	PENJET	42
JMIN=1	PENJET	43
C	PENJET	44
C LOCATING A PARTICLE IN 3-D SPACE...THEN FINDING WHETHER IT STRIKES	PENJET	45
C THE HOLE PROFILE. IF IT DOESN'T, THEN TIME IS INCREMENTED UNTIL	PENJET	46
C IMPACT OCCURS.	PENJET	47
C	PENJET	48
201 T(I)=T(I)+DELT	PENJET	49
205 X(I)=VEL(I)/AVST(I)+T(I)*COS(BETA(I))	PENJET	50
Y(I)=X(I)*TAN(BETA(I))	PENJET	51
Z(I)=VEL(I)*T(I)-ZSUBO+XL(I)+(DS(1)+DS(2)+DS(3))/3.	PENJET	52
IF(INO.EQ.2) GO TO 210	PENJET	53
IF(Z(I).LT.0.0) GO TO 201	PENJET	54
IF(T.EQ.1 .AND. T(I).LT.BREAK(I))GOTO 202	PENJET	55
IF(T.EQ.1)GO TO 210	PENJET	56
KLESS=4-1	PENJET	57
DO 206 J=JMIN,KLESS	PENJET	58
IF(PEN(J).LT..01)GO TO 206	PENJET	59
V(I,1)=X(I)-X(J)	PENJET	60
V(I,2)=Y(I)-Y(J)	PENJET	61
V(I,3)=Z(I)-Z(J)	PENJET	62
VM(I)=SQRT(V(I,1)**2+V(I,2)**2+V(I,3)**2)	PENJET	63
VAN(I)=V(I,1)*GRAD(J,1)+V(I,2)*GRAD(J,2)+V(I,3)*GRAD(J,3)	PENJET	64
VAM(I)=ABS(VAN(I))	PENJET	65
VPM(I)=SQRT(ABS(VM(I)**2-VAN(I)**2))	PENJET	66
IF(VPM(I)*DIA(I)/2.1**2/RADIUS(J)**2+VAN(I)**2/PEN(J)**2.GT.1.0)	PENJET	67
I GO TO 206	PENJET	68
IF(INO.EQ.1)GO TO 210	PENJET	69
JMIN=J	PENJET	70
GO TO 201	PENJET	71
206 CONTINUE	PENJET	72
207 NO=1	PENJET	73
IF(JMTN.EQ.KLESS .AND. Z(I).LE..1) GO TO 208	PENJET	74
T(I)=T(I)-DELT/6.0	PENJET	75
GO TO 205	PENJET	76
208 T(I)=(ZSUBO-XL(I)-(DS(1)+DS(2)+DS(3))/3.)/VEL(I)	PENJET	77
NO=2	PENJET	78
GO TO 205	PENJET	79
C	PENJET	80
C FINDING THE EFFECTS OF THE PARTICLE PENETRATION	PENJET	81
C	PENJET	82
210 IF(I.NE.1 .AND. NO.NE.2)GO TO 202	PENJET	83
GRAD(1,1)=0.0	PENJET	84
GRAD(1,2)=0.0	PENJET	85
GRAD(1,3)=1.0	PENJET	86
GO TO 281	PENJET	87
C	PENJET	88
C IF A PARTICLE IS NOT IN THE CONTINUOUS MODE, THE AXIS OF	PENJET	89
C PENETRATION IS EVALUATED	PENJET	90

C	202 IF(I(I).LT.BREAK(I))NCC(I)=1	PENJET	91
	IF(I.EQ.1 .AND. NCC(I).EQ.1)GOTO 1999	PENJET	92
	IF(I(I).LT.BREAK(I))XL(I)=XL(I)*T(I)/BREAK(I)	PENJET	93
	IM1=I-1	PENJET	94
	IF(I(I).LT.BREAK(I) .AND. Z(I).GT.Z(IM1))GOTO 1999	PENJET	95
	DO 204 LJ=1,3	PENJET	96
	VA(I,LJ)=GRAD(JMIN,LJ)*VAN(I)	PENJET	97
	VP(I,LJ)=V(I,LJ)-VA(I,LJ)	PENJET	98
	AR(I,LJ)=VA(I,LJ)/VAN(I)	PENJET	99
	204 PR(I,LJ)=VP(I,LJ)/VPM(I)	PENJET	100
	RGRAD(I,1)=VPM(I)*2./RADIUS(JMIN)**2	PENJET	101
	RGRAD(I,2)=0.0	PENJET	102
	RGRAD(I,3)=VAN(I)*2./PEN(JMIN)**2	PENJET	103
	RGM(I)=SQRT(RGRAD(I,1)**2+RGRAD(I,3)**2)	PENJET	104
	RGPDI(I,1)=RGRAD(I,1)/RGM(I)	PENJET	105
	RGRAD(I,3)=RGRAD(I,3)/RGM(I)	PENJET	106
	DO 214 LJ=1,3	PENJET	107
	214 GRAD(I,LJ)=RGRAD(I,LJ)*PR(I,LJ)+RGRAD(I,3)*AR(I,LJ)	PENJET	108
	GM=SQRT(GRAD(I,1)**2+GRAD(I,2)**2+GRAD(I,3)**2)	PENJET	109
	DO 219 LJ=1,3	PENJET	110
	219 GRAD(I,LJ)=GRAD(I,LJ)/GM	PENJET	111
	GOTO 251	PENJET	112
C		PENJET	113
C	CONTINUOUS PENETRATION AXIS IS ALONG PARTICLE FLIGHT AXIS.	PENJET	114
C		PENJET	115
C	1999 VECT=SQRT(X(I)**2+Y(I)**2+(Z(I)+ZSUBO-(DS(1)+DS(2)+DS(3))/3.)*2)	PENJET	116
	GRAD(I,1)=X(I)/VECT	PENJET	117
	GRAD(I,2)=Y(I)/VECT	PENJET	118
	GRAD(I,3)=(Z(I)+ZSUBO-(DS(1)+DS(2)+DS(3))/3.)/VECT	PENJET	119
	251 PENMAX(KZSUM)=XL(I)/GAMMA+PENMAX(KZSUM)	PENJET	120
	IF(I.EQ.1 .OR. NO.EQ.2 .OR. NCC(I).EQ.1)GOTO 262	PENJET	121
C		PENJET	122
C	EFFECTIVE LENGTH AND EFFECTIVE VELOCITY CALCULATIONS.	PENJET	123
C		PENJET	124
C	XL(I)=GRAD(I,3)*XL(I)+SQRT(GRAD(I,1)**2+GRAD(I,2)**2)	PENJET	125
	*DTA(I)	PENJET	126
	1 VEL(I)=GRAD(I,3)*VEL(I)	PENJET	127
C		PENJET	128
C	PENETRATION CALCULATIONS AS A FUNCTION OF EFFECTIVE VELOCITY.	PENJET	129
C		PENJET	130
C	IF(VEL(I).GT. .8) GO TO 260	PENJET	131
	PEN(I)=0.0	PENJET	132
	RADIUS(I)=0.0	PENJET	133
	U(I)=0.0	PENJET	134
	HVOL(I)=0.0	PENJET	135
	GO TO 209	PENJET	136
	260 IF(VEL(I).GT.1.4) GO TO 261	PENJET	137
	U(I)=0.0	PENJET	138
	HVOL(I)=XKE(I)/ECONST	PENJET	139
	PEN(I)=(VEL(I)-.8)**2/.72*XL(I)	PENJET	140
	RADIUS(I)=SQRT(3.*HVOL(I) / (2.*3.14159*PEN(I)))	PENJET	141
	GO TO 209	PENJET	142
	261 IF(VEL(I).GT.2.1)GO TO 262	PENJET	143
	U(I)=0.0	PENJET	144
	HVOL(I)=XKE(I)/ECONST	PENJET	145
	PEN(I)=(-(VEL(I)-2.)*2/.72*1.)*XL(I)	PENJET	146
	RADIUS(I)=SQRT(3.*HVOL(I) / (2.*3.14159*PEN(I)))	PENJET	147
	GO TO 209	PENJET	148
	262 RAD=SQRT(VEL(I)**2+RHOJ*RHOT-2.*(RHOJ-RHOT)*(SIGJ-SIGT)/1.E6)	PENJET	149
	U(I)=(RHOJ*VEL(I)-RAD)/(RHOJ-RHOT)	PENJET	150
	PEN(I)=(RHOJ*VEL(I)*XL(I)-XL(I)*RAD)/(RAD-RHOT*VEL(I))	PENJET	151
	HVOL(I)=XKE(I)/ECONST	PENJET	152
	RADIUS(I)=SQRT(3.*HVOL(I) / (2.*3.14159*PEN(I)))	PENJET	153
		PENJET	154

209 CONTINUE	PENJET	155
279 WRITE(6,211)	PENJET	156
211 FORMAT(1H1,10X,'PARTICLE',5X,'HOLE',6X,'IMPACT',3X,'PENETRATION ',	PENJET	157
1'PENETRATION',6X,'STATUS',12X,'NUMBER',5X,'VOLUME',6X,'TIME',17X,	PENJET	158
2'VELOCITY',7,23X,'(CU.MM) (MICROSEC)',4X,'(MM)',5X,'(MM/)',	PENJET	159
3'MICROSEC)',/)	PENJET	160
DO 213 I=1,NPART	PENJET	161
IF(INCC(I).EQ.1)GOTO 281	PENJET	162
WRITE(6,212)I,HVOL(I),T(I),PEN(I),U(I),MPC,MPC,MPC,MPC	PENJET	163
GO TO 213	PENJET	164
281 WRITE(6,212)I,HVOL(I),T(I),PEN(I),U(I),MPC,MPC,MPC,MPC	PENJET	165
212 FORMAT(14X,12,7X,F6.0,4X,F5.1,7X,F4.1,8X,F5.3,7X,3A4)	PENJET	166
213 CONTINUE	PENJET	167
PMAX=0.0	PENJET	168
DO 215 I=1,NPART	PENJET	169
IF(Z(I).LT.PMAX) GO TO 215	PENJET	170
PMAX=Z(I)	PENJET	171
DPHAX=PEN(I)*GRAD(I,3)	PENJET	172
215 CONTINUE	PENJET	173
PENSUM=PMAX+DPHAX	PENJET	174
PENT(KZSUM)=PENSUM	PENJET	175
WRITE(6,216)PENSUM	PENJET	176
216 FORMAT(1H0,10X,'THE TOTAL PENETRATION PREDICTED BY THIS ',	PENJET	177
1'PROGRAM IS ',F6.1,' MM')	PENJET	178
PRINT 257, ZSUBO	PENJET	179
257 FORMAT(1H0,10X,'THESE VALUES WERE CALCULATED AT A STANDOFF ',	PENJET	180
1'OF ',F6.0,' MM')	PENJET	181
WRITE(6,220)	PENJET	182
220 FORMAT(1H1,14X,'X,Y, AND Z COORDINATES PENETRATION MAX CRA',	PENJET	183
1'TER RADIUS AXIAL DIRECTION OF PENETRATION')	PENJET	184
DO 218 I=1,NPART	PENJET	185
WRITE(6,217)I,X(I),Y(I),Z(I),PEN(I),RADIUS(I),GRAD(I,1),	PENJET	186
1GRAD(I,2),GRAD(I,3)	PENJET	187
217 FORMAT(11X,12,9F10.2,3F10.4)	PENJET	188
218 CONTINUE	PENJET	189
DO 263 I=1,NPART	PENJET	190
VEL(I)=RVEL(I)	PENJET	191
KWAL(I)=0	PENJET	192
263 XL(I)=RL(I)	PENJET	193
IF(ZSUBO.LT.ZMAX) GO TO 250	PENJET	194
WRITE(6,290)	PENJET	195
290 FORMAT(1H1,' VIRTUAL ORIGIN ZO(MM) PENETRATION(MM) MAX ',	PENJET	196
1'PENETRATION(MM)',/)	PENJET	197
DO 289 I=1,KZSUM	PENJET	198
289 WRITE(6,292)ZSUB(I),PENT(I),PENMAX(I)	PENJET	199
292 FORMAT(10X,F5.0,T32,F7.2,T55,F7.2)	PENJET	200
515 CONTINUE	PENJET	201
505 FORMAT(1H1)	BRL2330	237
DO 510 JN= 1,NPART	BRL2330	238
VOL(JN)=0.	BRL2330	239
XL(JN)=0.	BRL2330	240
510 DIA(JN)=0.	BRL2330	241
500 CONTINUE	BRL2330	242
STOP	BRL2330	243
END	BRL2330	244

"JETBREAK"

BREAK(N) - breakup time of particle N in μ sec

BROK $\left. \begin{matrix} 1 \\ 2 \\ 3 \end{matrix} \right\}$ - observed number of jet particles already particulated as of
 $\begin{matrix} 1 \\ 2 \\ 3 \end{matrix}$
 exposure $\begin{matrix} 1 \\ 2 \\ 3 \end{matrix}$

FLASH $\left. \begin{matrix} 1 \\ 2 \\ 3 \end{matrix} \right\}$ - time after detonation of round for exposure of film $\begin{matrix} 1 \\ 2 \\ 3 \end{matrix}$

PBAR -
 PBYPZ -
 PBYTZ -
 TAVG -
 $\left. \begin{matrix} - \\ - \\ - \end{matrix} \right\}$ statistical variables used in calculation of breakup time

TDLAY - time delay between detonation of round and formation of jet

Z1 - time between breakup of adjacent jet particles

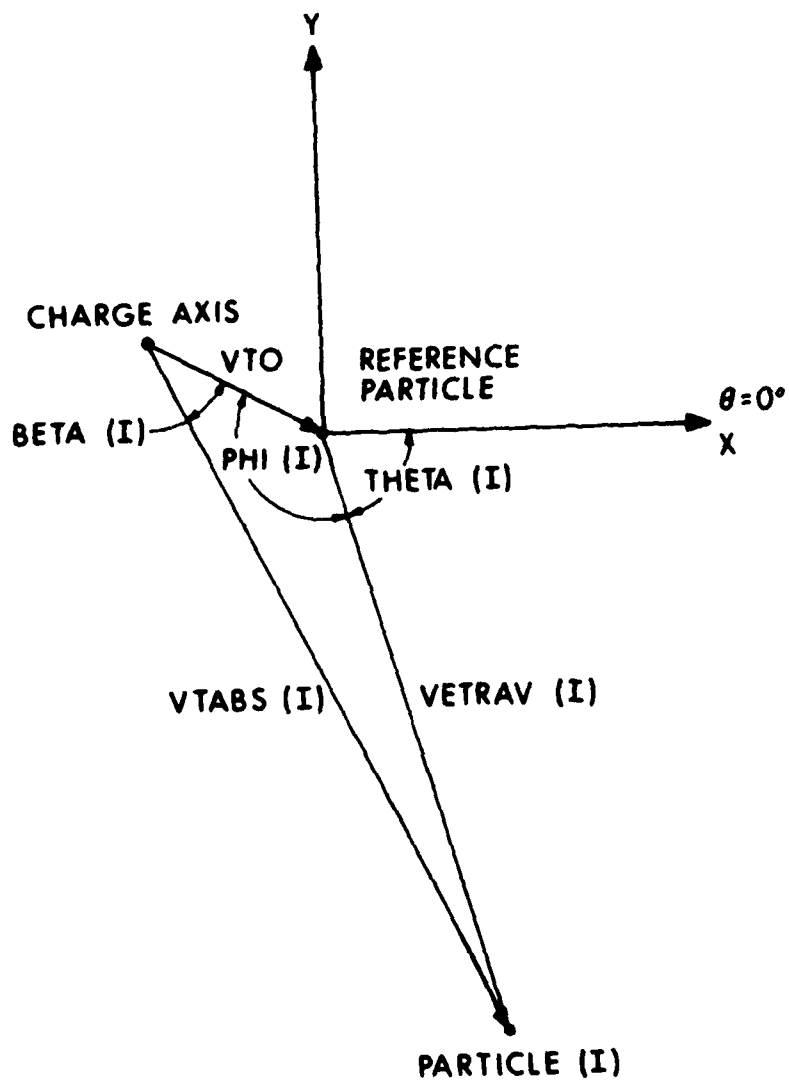


Figure A1. A Diagrammatic Sketch Identifying Some DRIFT Variables

"DRIFT"

- AVST (I) - the axial velocity divided by the transverse velocity for particle I
- AZ (1, I) }
 AZ (2, I) } - deviance of particle I with respect to the reference on
 AZ (3, I) } exposure 1
 exposure 2
 exposure 3
- BETA (I) - absolute angle formed between drifting reference and particle I
- DS (1) }
 DS (2) } - virtual origin calculations based on exposure 1
 DS (3) } exposure 2
 exposure 3
- EPS (1, I) }
 EPS (2, I) }
 EPS (3, I) } - estimate for sensitivity of transverse velocity V with respect to
 EPS (4, I) } small changes in drift angle θ (of the form $\frac{dV}{d\theta} \Delta\theta$) using all
 EPS (5, I) } combinations of drift velocities and drift angles calculated from
 the 3 radiographs
- EPSUM - $\sum_{K=1}^6 \frac{1}{EPS(K, I)}$ a constant used to scale the EPS values to their proper magnitude
- KD 12 }
 KD 23 } - a constant relating time delay between exposures 1 2
 KD 13 } drift velocity calculations 1 3
- MAXDEX - the particle # that satisfies TMAX
- MINDEX - the particle # that satisfies TMIN
- PHI (I) - the supplementary of the relative drift angle between particle I and the reference
- R (2) }
 R (1) } - correlation coefficient for V_A vs. V_T plot for reference drift
 angle presently being analyzed.
 previously
- S1 (I) }
 S2 (I) } - axial distance from charge base to particle I on exposure 1
 S3 (I) } exposure 2
 exposure 3

T1 (I) }	- estimated existence time from creation of particle I to exposure		
T2 (I) }		1	
T3 (I) }	of film 2.		
		3	
THETA0	- initial guess at direction of reference drift		
TMAX	- the angle of the particle which is most clockwise/counterclockwise		
TMIN	in the relative angular span.		
TH12 }		1	2
TH23 }	- relative drift angle calculated from governing equations	2	and 3
TH13 }		1	3
THV	- the numerical average of TH12, TH23, and TH13		
THETA (I)	- the relative drift angle of particle I with respect to arbitrary origin (chosen parallel to film 2) (degrees).		
THOR	- the current reference drift angle being tested		
VETR12 }		1	2
VETR23 }	- relative transverse velocity calculated from equations	2	and 3
VETR13 }		1	3
VETRAV (I)	- relative transverse velocity of particle I based on weighted values of VETR12, VETR23, and VETR13 (m/s)		
VTABS (I)	- absolute transverse velocity of particle I (m/s)		
VTO	- transverse velocity of particle reference (m/s)		
VXO	- x component of drift velocity for the reference		
VYO	- y		
VX (I)	- x component of relative drift velocity for particle I		
VY (I)	- y		
WT (I)	- the percentage weight that an individual drift velocity calculation contributes to the overall average.		
ABAR			
ABYAS			
AS			
CO			
CO1			
CO2			
TBAR			
TBYAS			
TBYTS			
TS			
XINT			
	statistical variables used in V_A vs V_T regression		

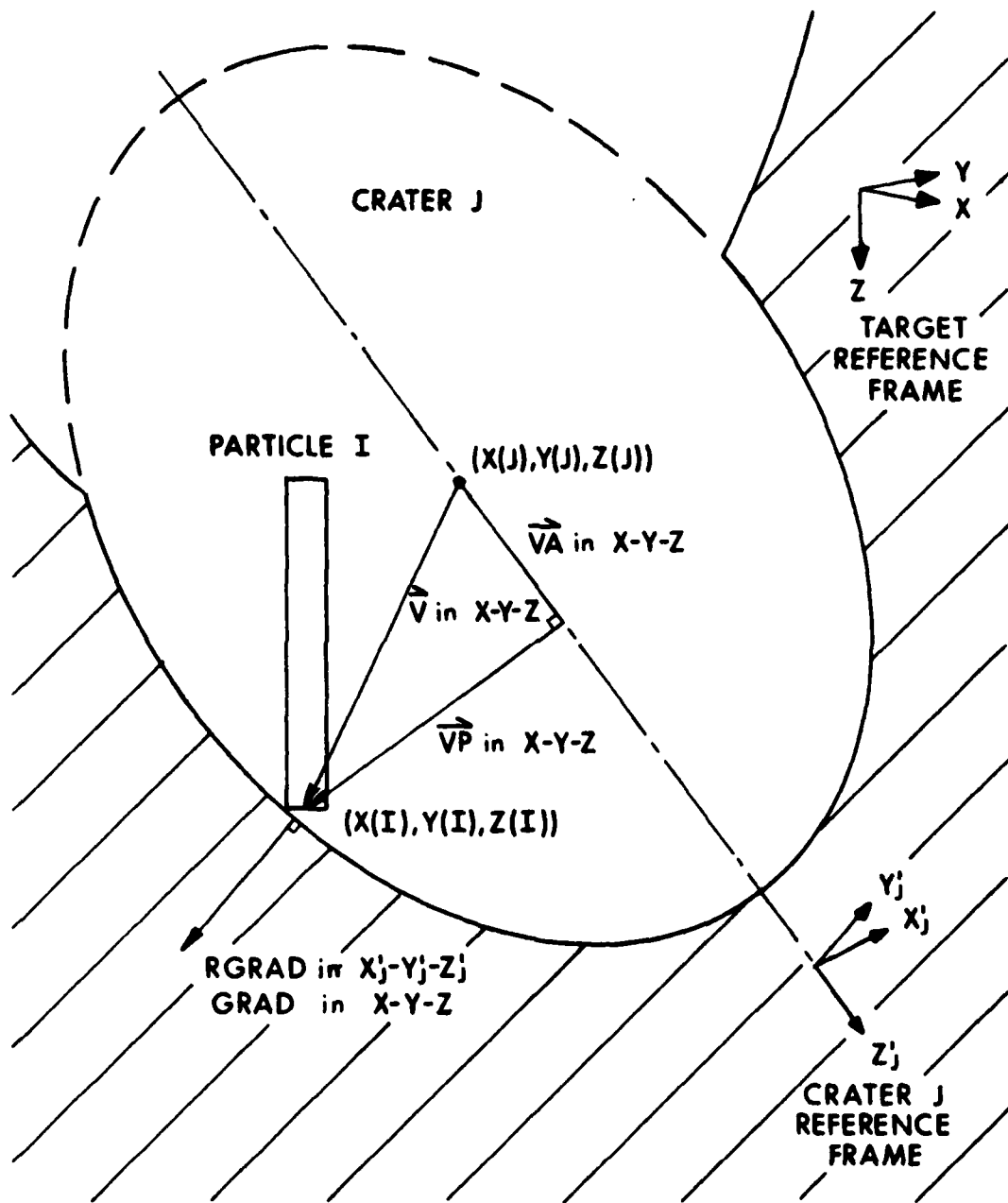


Figure A2: A Diagrammatic Sketch Identifying Some PENJET Variables.

"PENJET"

- T (I) - impact time of particle I, taken from virtual origin
- X (I) }
Y (I) } - spacial coordinates for top of particle I at any given time
Z (I) } (and eventually upon impact) X = 0, Y = 0 denote where charge axis intersects target plate; Z = 0 is target surface.
- DELT - time increment (varied so that particle travels 3 mm/timestep)
- U (I) - penetration velocity of particle I (calculated only in the hypervelocity regime e.g. above 2 mm/ μ sec)
- RADIUS (I) - crater radius of particle I
- PEN (I) - penetration of particle I
- HVOL (I) - hole volume created by particle I
- V (I, 1) }
V (I, 2) }
V (I, 3) } $\begin{matrix} x \\ y \\ z \end{matrix}$ the y components of a vector drawn from the center of crater J (through which particle I is instantaneously passing) to the tip of particle I
- VM (I) - magnitude of vector \vec{V}
- VA (I, 1) }
VA (I, 2) }
VA (I, 3) } $\begin{matrix} x \\ y \\ z \end{matrix}$ the y components of vector \vec{VA} which is in itself the component of vector \vec{V} pointing axially down crater J (the crater through which particle I is passing)
- VAM (I) - magnitude of vector \vec{VA}
- VP (I, 1) }
VP (I, 2) }
VP (I, 3) } $\begin{matrix} x \\ y \\ z \end{matrix}$ the y components of vector \vec{VP} which is in itself the component of vector \vec{V} pointing perpendicular to the penetration axis of crater J (through which particle I is passing).
Note that $\vec{VA} + \vec{VP} = \vec{V}$
- VPM (I) - magnitude of vector \vec{VP}
- RGRAD (I, 1) }
RGRAD (I, 2) }
RGRAD (I, 3) } $\begin{matrix} x \\ y \\ z \end{matrix}$ the y components of a vector that is perpendicular to crater J at the point of impact of particle I, but in the coordinate system of crater J (e.g., the Z' axis points down the axis of crater J)
- RGM (I) - magnitude of vector \vec{RGRAD}

GRAD (I, 1) } x
 GRAD (I, 2) } the y components of a vector that is perpendicular to crater J
 GRAD (I, 3) } z
 at the point of impact of particle I, and in laboratory coordinate system (e.g., the z axis points perpendicular to the target surface); GRAD becomes the axis of penetration for crater I according to refracted particle wave theory

GM (I) - magnitude of vector $\overrightarrow{\text{GRAD}}$

AR (I, 1) }
 AR (I, 2) } a unit vector version of $\overrightarrow{\text{VA}}$
 AR (I, 3) }

PR (I, 1) }
 PR (I, 2) } a unit vector version of $\overrightarrow{\text{VP}}$
 PR (I, 3) }

NCC (I) - equals 1 when particle I is particulate; equals 0 when I is continuous

ECONST - energy constant used for calculating crater volume (J/mm^3)

ZSUBO - standoff (mm) from virtual origin to target surface

RHOJ - jet density ρ_j (g/mm^3)

RHOT - target density ρ_t (g/mm^3)

SIGJ - jet strength (MPa)

SIGT - target strength (MPa)

GAMMA - $(\rho_t / \rho_j)^{1/2}$

JMIN - the last crater number that particle I has been known to successfully have passed through

PENSUM - the penetration resulting from PENJET calculations

NQ = 0 - particle I is travelling down hole profile

NQ = 1 - particle I has just struck target; t is decremented until precise impact time is revealed

NQ = 2 - particle I has drifted so far off course as to preclude entry into any part of the hole profile; it therefore strikes the target surface

KZSUM - number of standoffs at which PENJET has just analyzed a drift distribution

PENMAX (K) - the greatest possible penetration at standoff #K

ZSUB (K) - an array of the standoffs at which PENJET has just been run
PENT (K) - penetration predicted by PENJET at standoff #K

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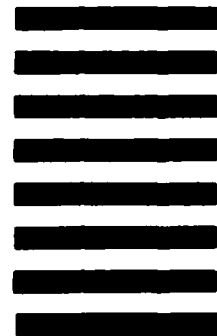


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